

An energy conserving, upwinded compatible finite element discretisation for the dry compressible Euler equations

Golo Wimmer, Colin Cotter, Werner Bauer

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Outline

1. Background

- Compatible finite elements
- Hamiltonian framework
- Upwind stabilization methods

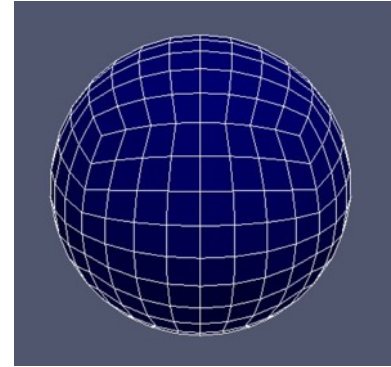
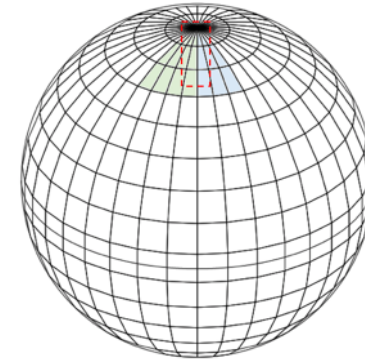
2. Hamiltonian framework including upwinding for cFEM

- Results for compressible Euler equations
- Computational efficiency
- Interpretation

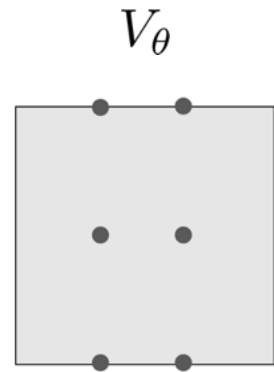
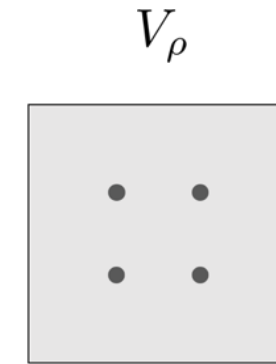
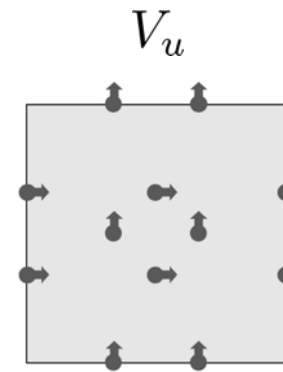
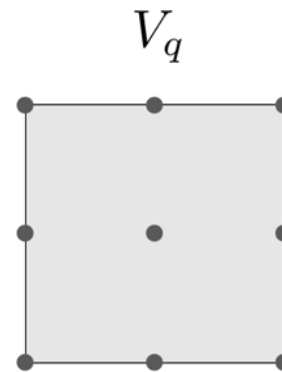
Compatible finite elements

- Used in UK Met Office's next dynamical core GungHo
 - FEM extension of Arakawa C Grid
 - Allows for general grids

- Map FE spaces via differential operators



$$\begin{array}{ccccc}
 H^1(\Omega) & \xrightarrow{\nabla \times} & H(\text{div}; \Omega) & \xrightarrow{\nabla \cdot} & L^2(\Omega) \\
 \downarrow \pi^0 & & \downarrow \pi^1 & & \downarrow \pi^2 \\
 V^0(\Omega) & \xrightarrow{\nabla \times} & V^1(\Omega) & \xrightarrow{\nabla \cdot} & V^2(\Omega)
 \end{array}$$



Vertical slice elements

CP staggering

Energy conservation in NWP

$$\left. \begin{aligned} \frac{D\mathbf{u}}{Dt} + 2\boldsymbol{\Omega} \times \mathbf{u} + c_p\theta\nabla\pi &= -g\mathbf{k} \\ \frac{D\rho}{Dt} + \rho\nabla \cdot \mathbf{u} &= 0 \\ \frac{D\theta}{Dt} &= 0 \end{aligned} \right\} \begin{aligned} E &= K + P + I \\ \frac{dE}{dt} &= \frac{dK}{dt} + \frac{dP}{dt} + \frac{dI}{dt} = 0 \end{aligned}$$

- Failure to account for a numerically closed energy budget may lead to net energy biases in climate models
 - Partly due to missing reinjection $K \rightarrow I$

Energy conservation in NWP

$$\frac{D\mathbf{u}}{Dt} + 2\boldsymbol{\Omega} \times \mathbf{u} + c_p\theta\nabla\pi = -g\mathbf{k}$$

$$\frac{D\rho}{Dt} + \rho\nabla \cdot \mathbf{u} = 0$$

$$\frac{D\theta}{Dt} = 0$$

$$E = K + P + I$$

$$\frac{dE}{dt} = \frac{dK}{dt} + \frac{dP}{dt} + \frac{dI}{dt} = 0$$

$$\frac{dE}{dt} = \frac{\delta E}{\delta \mathbf{u}} \frac{d\mathbf{u}}{dt} + \frac{\delta E}{\delta \rho} \frac{d\rho}{dt} + \frac{\delta E}{\delta \theta} \frac{d\theta}{dt}$$



Discretized equations

Hamiltonian framework

■ Hamiltonian representing total energy $H(\mathbf{u}, \rho, \theta) = \int_{\Omega} \left(\frac{\rho}{2} |\mathbf{u}|^2 + g\rho z + c_v \rho \theta \pi \right) dx$

■ Poisson bracket $\{F, H\} = -\left\langle \frac{\delta F}{\delta \mathbf{u}}, \boldsymbol{\omega} \times \frac{\delta H}{\delta \mathbf{u}} \right\rangle + \left\langle \nabla \cdot \frac{\delta F}{\delta \mathbf{u}}, \frac{\delta H}{\delta \rho} \right\rangle + \left\langle \frac{\delta H}{\delta \theta}, \frac{1}{\rho} \frac{\delta F}{\delta \mathbf{u}} \cdot \nabla \theta \right\rangle$

- Bilinear form
- Antisymmetric $-\left\langle \nabla \cdot \frac{\delta H}{\delta \mathbf{u}}, \frac{\delta F}{\delta \rho} \right\rangle - \left\langle \frac{\delta F}{\delta \theta}, \frac{1}{\rho} \frac{\delta H}{\delta \mathbf{u}} \cdot \nabla \theta \right\rangle$

$$\frac{\delta H}{\delta \mathbf{u}} = \rho \mathbf{u}$$

$$\frac{\delta H}{\delta \rho} = \frac{1}{2} |\mathbf{u}|^2 + gz + c_p \theta \pi$$

$$\frac{\delta H}{\delta \theta} = c_p \rho \pi$$

$$\pi = \left(\frac{R}{p_0} \rho \theta \right)^{\frac{1-\kappa}{\kappa}}$$

$$\boldsymbol{\omega} = \frac{1}{\rho} (\nabla \times \mathbf{u} + \boldsymbol{\Omega})$$

Hamiltonian framework

■ Poisson system $\frac{dF}{dt} = \{F, H\}$

■ Poisson bracket
$$\{F, H\} = -\left\langle \frac{\delta F}{\delta \mathbf{u}}, \omega \times \frac{\delta H}{\delta \mathbf{u}} \right\rangle + \left\langle \nabla \cdot \frac{\delta F}{\delta \mathbf{u}}, \frac{\delta H}{\delta \rho} \right\rangle + \left\langle \frac{\delta H}{\delta \theta}, \frac{1}{\rho} \frac{\delta F}{\delta \mathbf{u}} \cdot \nabla \theta \right\rangle$$

$-\left\langle \nabla \cdot \frac{\delta H}{\delta \mathbf{u}}, \frac{\delta F}{\delta \rho} \right\rangle$
 $-\left\langle \frac{\delta F}{\delta \theta}, \frac{1}{\rho} \frac{\delta H}{\delta \mathbf{u}} \cdot \nabla \theta \right\rangle$

$F = \langle \phi, \rho \rangle \Rightarrow \frac{\delta F}{\delta \rho} = \phi, \quad \left\langle \phi, \frac{\partial \rho}{\partial t} \right\rangle = -\left\langle \phi, \nabla \cdot (\rho \mathbf{u}) \right\rangle$

$F = \langle \gamma, \theta \rangle \Rightarrow \frac{\delta F}{\delta \theta} = \gamma, \quad \left\langle \gamma, \frac{\partial \theta}{\partial t} \right\rangle = -\left\langle \gamma, \mathbf{u} \cdot \nabla \theta \right\rangle$

Hamiltonian framework

■ Poisson system $\frac{dF}{dt} = \{F, H\}$

■ Poisson bracket
$$\begin{aligned} \{F, H\} = & -\left\langle \frac{\delta F}{\delta \mathbf{u}}, \omega \times \frac{\delta H}{\delta \mathbf{u}} \right\rangle + \left\langle \nabla \cdot \frac{\delta F}{\delta \mathbf{u}}, \frac{\delta H}{\delta \rho} \right\rangle + \left\langle \frac{\delta H}{\delta \theta}, \frac{1}{\rho} \frac{\delta F}{\delta \mathbf{u}} \cdot \nabla \theta \right\rangle \\ & - \left\langle \nabla \cdot \frac{\delta H}{\delta \mathbf{u}}, \frac{\delta F}{\delta \rho} \right\rangle - \left\langle \frac{\delta F}{\delta \theta}, \frac{1}{\rho} \frac{\delta H}{\delta \mathbf{u}} \cdot \nabla \theta \right\rangle \end{aligned}$$

$$F = H \quad \Rightarrow \quad \frac{dH}{dt} = \{H, H\} = -\{H, H\} = 0$$

Poisson bracket – no upwinding

$$\frac{dF}{dt} = \{F, H\} \quad \{F, H\} = -\left\langle \frac{\delta F}{\delta \mathbf{u}}, \omega \times \frac{\delta H}{\delta \mathbf{u}} \right\rangle + \left\langle \nabla \cdot \frac{\delta F}{\delta \mathbf{u}}, \frac{\delta H}{\delta \rho} \right\rangle + \left\langle \frac{\delta H}{\delta \theta}, \frac{1}{\rho} \frac{\delta F}{\delta \mathbf{u}} \cdot \nabla \theta \right\rangle$$

$$-\left\langle \nabla \cdot \frac{\delta H}{\delta \mathbf{u}}, \frac{\delta F}{\delta \rho} \right\rangle - \left\langle \frac{\delta F}{\delta \theta}, \frac{1}{\rho} \frac{\delta H}{\delta \mathbf{u}} \cdot \nabla \theta \right\rangle$$

$$\frac{\delta H}{\delta \mathbf{u}} = \mathbf{F} = P_{V_{\mathbf{u}}}(\rho \mathbf{u})$$

$$\frac{\delta H}{\delta \rho} = P = P_{V_{\rho}}\left(\frac{1}{2}|\mathbf{u}|^2 + gz + c_p \theta \pi\right)$$

$$\frac{\delta H}{\delta \theta} = T = P_{V_{\theta}}(c_p \rho \pi)$$

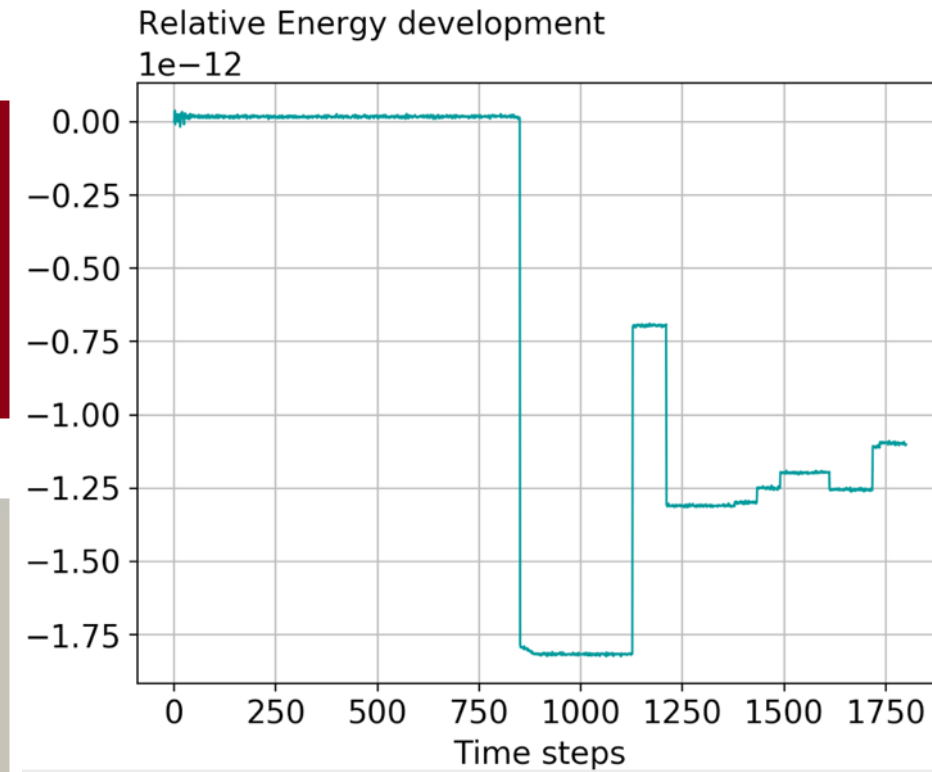
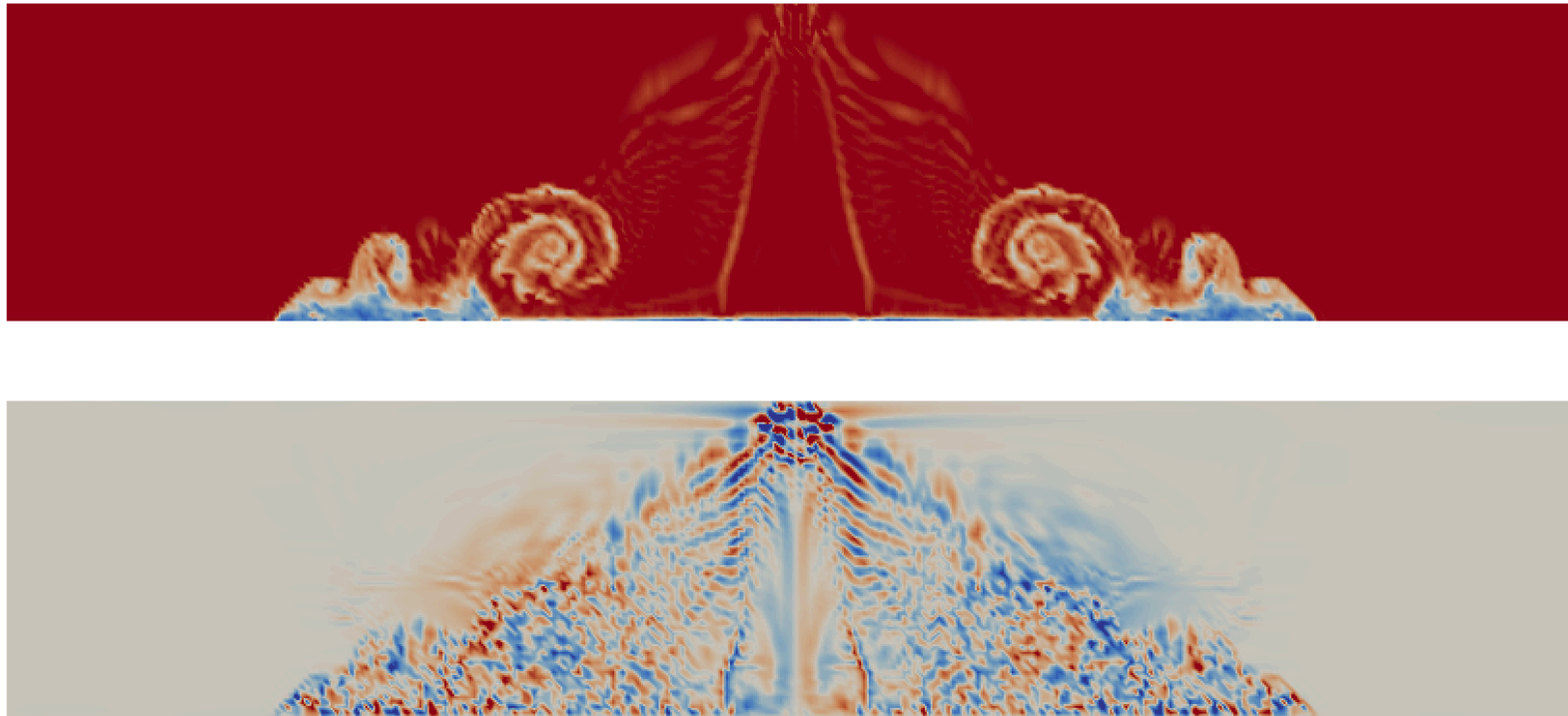
$$\langle \mathbf{w}, \mathbf{u}_t + \omega \times \mathbf{F} \rangle - \langle \nabla \cdot \mathbf{w}, P \rangle - \langle \mathbf{w}, T/\rho \nabla \theta \rangle = 0$$

$$\langle \phi, \rho_t + \nabla \cdot \mathbf{F} \rangle = 0$$

$$\langle \gamma, \theta_t + \mathbf{F}/\rho \cdot \nabla \theta \rangle = 0$$

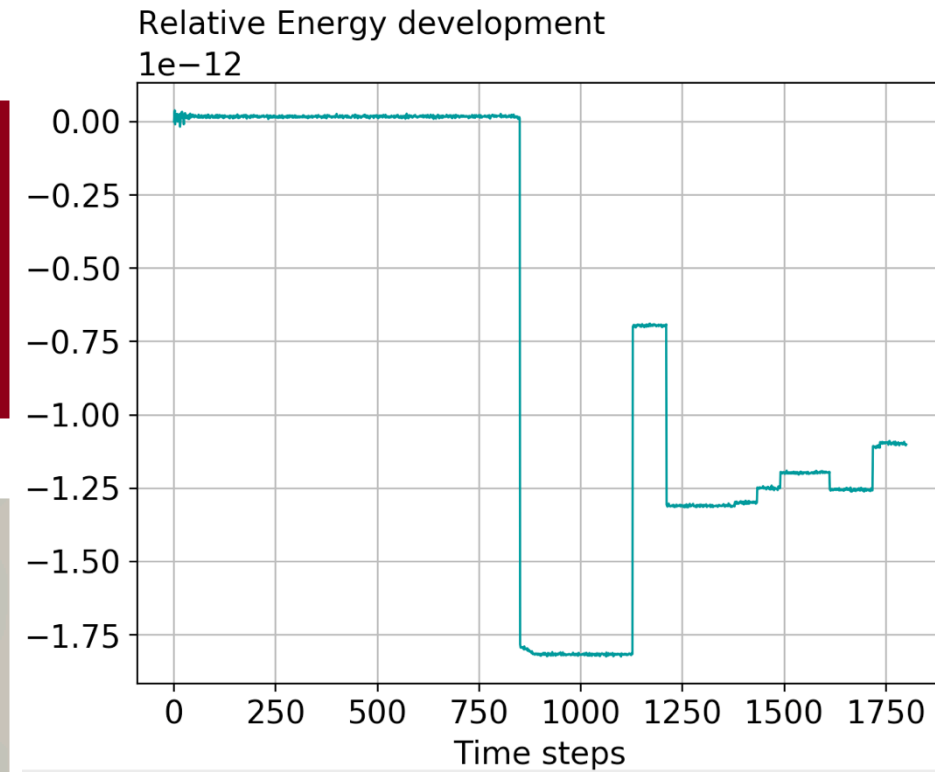
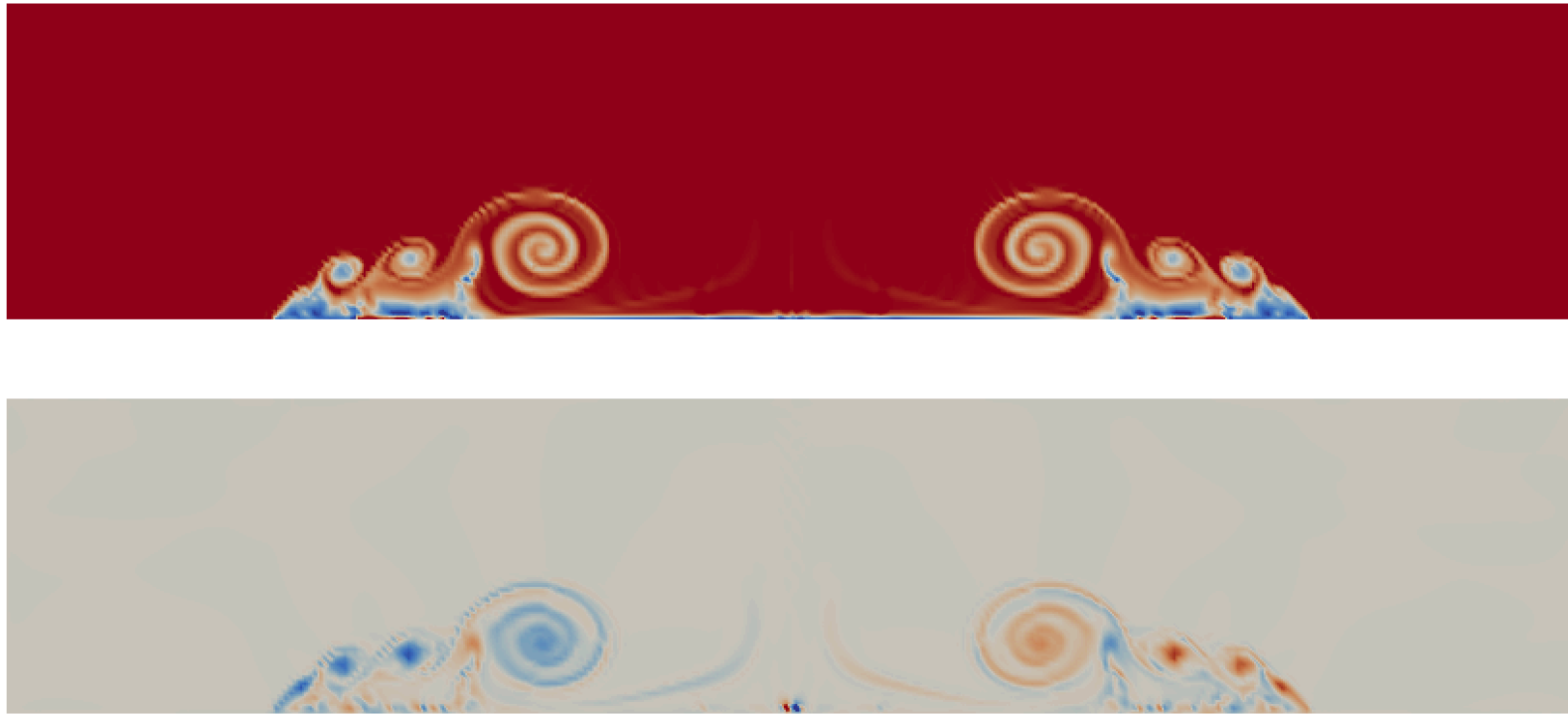
Test results – dry compressible Euler

- Falling cold bubble test case



Test results – dry compressible Euler

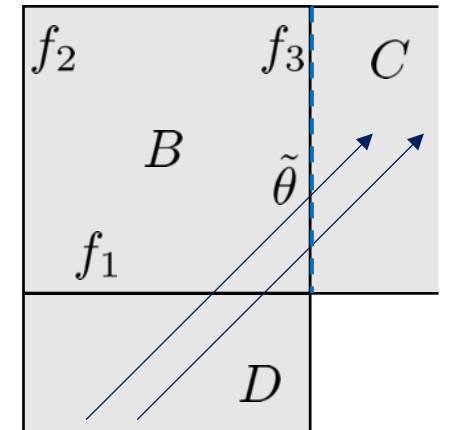
- Falling cold bubble test case



Upwind stabilization

- Different finite element spaces require different advection schemes
 - Horizontal DG upwinding** for potential temperature θ

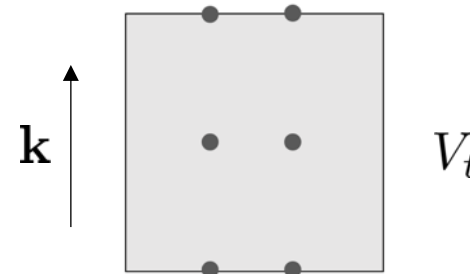
$$\begin{aligned} \langle \gamma, \mathbf{u} \cdot \nabla \theta \rangle &\longrightarrow -\langle \nabla \cdot (\gamma \mathbf{u}), \theta \rangle + \int_{\Gamma} [[\gamma \mathbf{u}]] \tilde{\theta} \, dS \\ &\longrightarrow \langle \gamma, \mathbf{u} \cdot \nabla \theta \rangle + \int_{\Gamma} ([[\gamma \mathbf{u}]] \tilde{\theta} - [[\gamma \mathbf{u} \theta]]) \, dS = L_{\mathbf{u}}(\theta; \gamma) \end{aligned}$$



- Vertical SUPG**

$$\gamma \longrightarrow \gamma + \tau(\mathbf{k} \cdot \mathbf{u})(\mathbf{k} \cdot \nabla \gamma) = \gamma + \tilde{\gamma}_{\mathbf{u}}$$

$$\langle \gamma + \tilde{\gamma}_{\mathbf{u}}, \theta_t \rangle = L_{\mathbf{u}}(\theta, \gamma + \tilde{\gamma}_{\mathbf{u}})$$



Potential temperature

- Move upwind part $\tilde{\gamma}_{\mathbf{u}}$ into bracket

$$\underline{\langle \gamma + \tilde{\gamma}_{\mathbf{u}}, \theta_t \rangle = L_{\mathbf{u}}(\theta, \gamma + \tilde{\gamma}_{\mathbf{u}})} \quad \longrightarrow \quad \underbrace{\langle \gamma, \theta_t \rangle}_{dF} = \underbrace{L_{\mathbf{u}}(\theta, \gamma + \tilde{\gamma}_{\mathbf{u}})}_{\{F, H\}} - \langle \tilde{\gamma}_{\mathbf{u}}, \theta_t \rangle$$

$$\frac{dF}{dt} = \{F, H\} \quad \text{for} \quad F = \langle \gamma, \theta \rangle$$

$$+\left\langle \frac{\delta H}{\delta \theta}, \frac{1}{\rho} \frac{\delta F}{\delta \mathbf{u}} \cdot \nabla \theta \right\rangle$$

$$+ L_{\mathbb{U}(\rho, \frac{\delta F}{\delta \mathbf{u}})}^{\theta} \left(\theta; \frac{\delta H}{\delta \theta} + \left(\frac{\widetilde{\delta H}}{\delta \theta} \right)_{\mathbf{u}} \right) + \left\langle \left(\frac{\widetilde{\delta H}}{\delta \theta} \right)_{\mathbf{u}}, \theta_t \right\rangle$$

$$-\left\langle \frac{\delta F}{\delta \theta}, \frac{1}{\rho} \frac{\delta H}{\delta \mathbf{u}} \cdot \nabla \theta \right\rangle$$

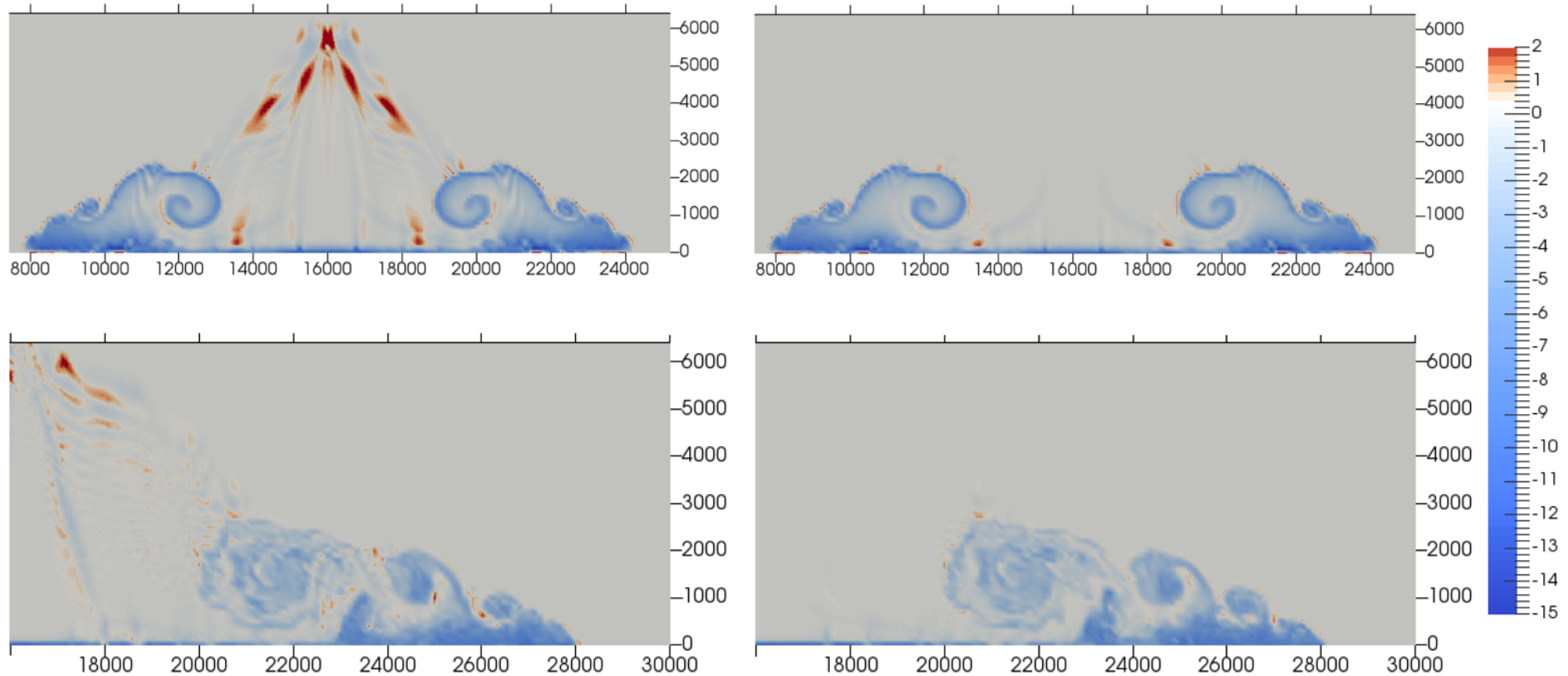
—————
DG/SUPG

$$- L_{\mathbb{U}(\rho, \frac{\delta H}{\delta \mathbf{u}})}^{\theta} \left(\theta; \frac{\delta F}{\delta \theta} + \left(\frac{\widetilde{\delta F}}{\delta \theta} \right)_{\mathbf{u}} \right) - \left\langle \left(\frac{\widetilde{\delta F}}{\delta \theta} \right)_{\mathbf{u}}, \theta_t \right\rangle$$

Results

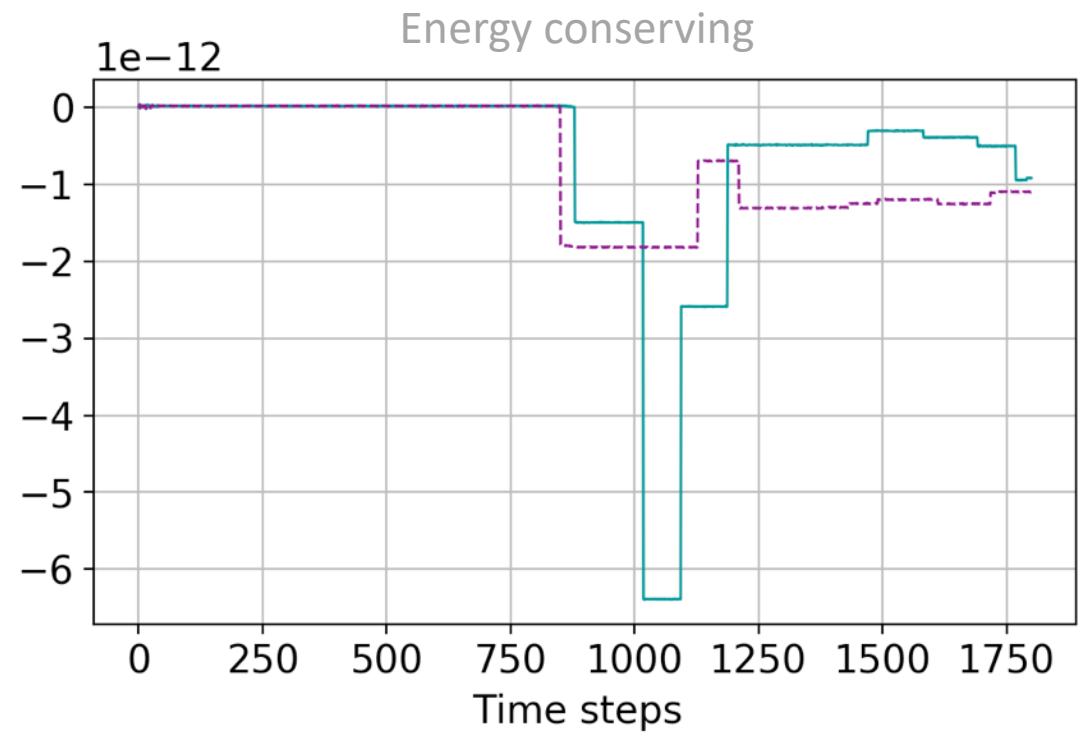
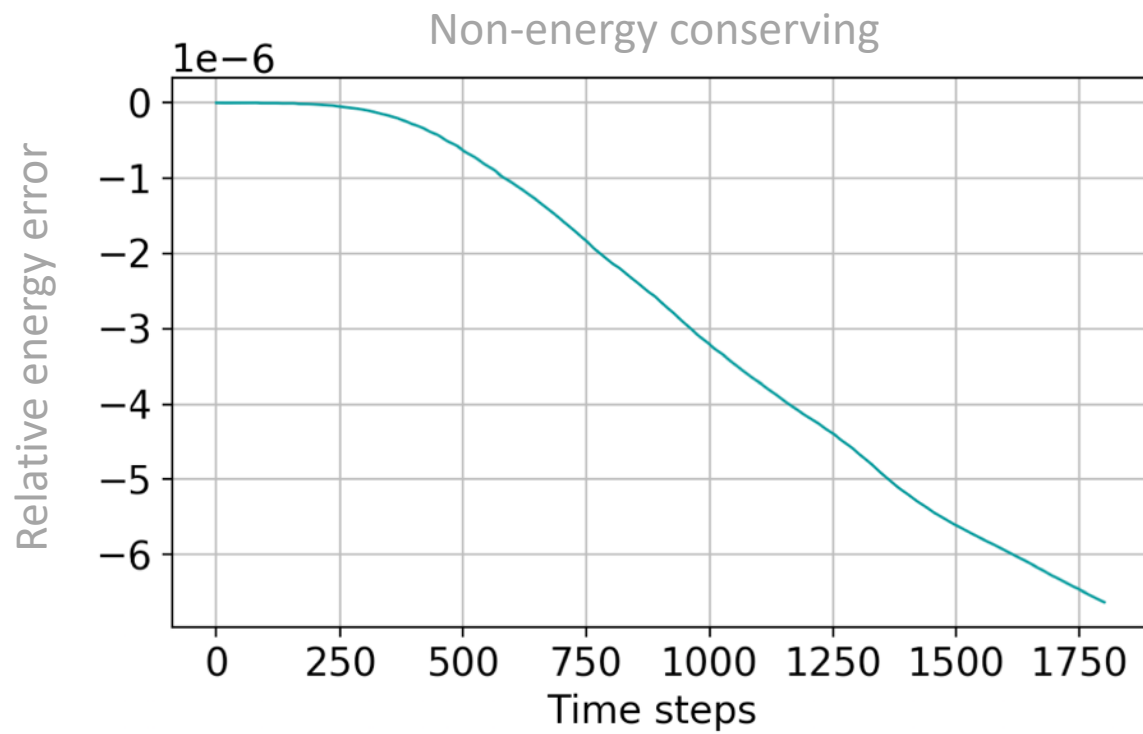
- Falling cold air bubble

Potential temperature



Results

- Falling cold air bubble



Computational cost

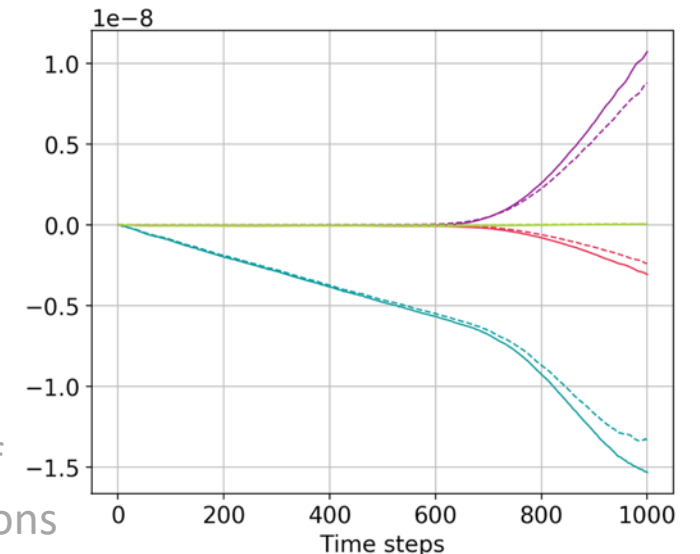
- Revert to approximately energy conserving time discretization

$$\overline{\frac{\delta H}{\delta \mathbf{u}}} = \frac{1}{3} P_{V_u} \left(\rho^n \mathbf{u}^n + \frac{1}{2} \rho^{n+1} \mathbf{u}^n + \frac{1}{2} \rho^n \mathbf{u}^{n+1} + \rho^{n+1} \mathbf{u}^{n+1} \right) \quad \longrightarrow \quad \overline{\frac{\delta H'}{\delta \mathbf{u}}} = P_{V_u} (\bar{\rho} \bar{\mathbf{u}})$$

- Can rearrange Poisson bracket swapping upwind and advecting velocities
- Remaining projections:

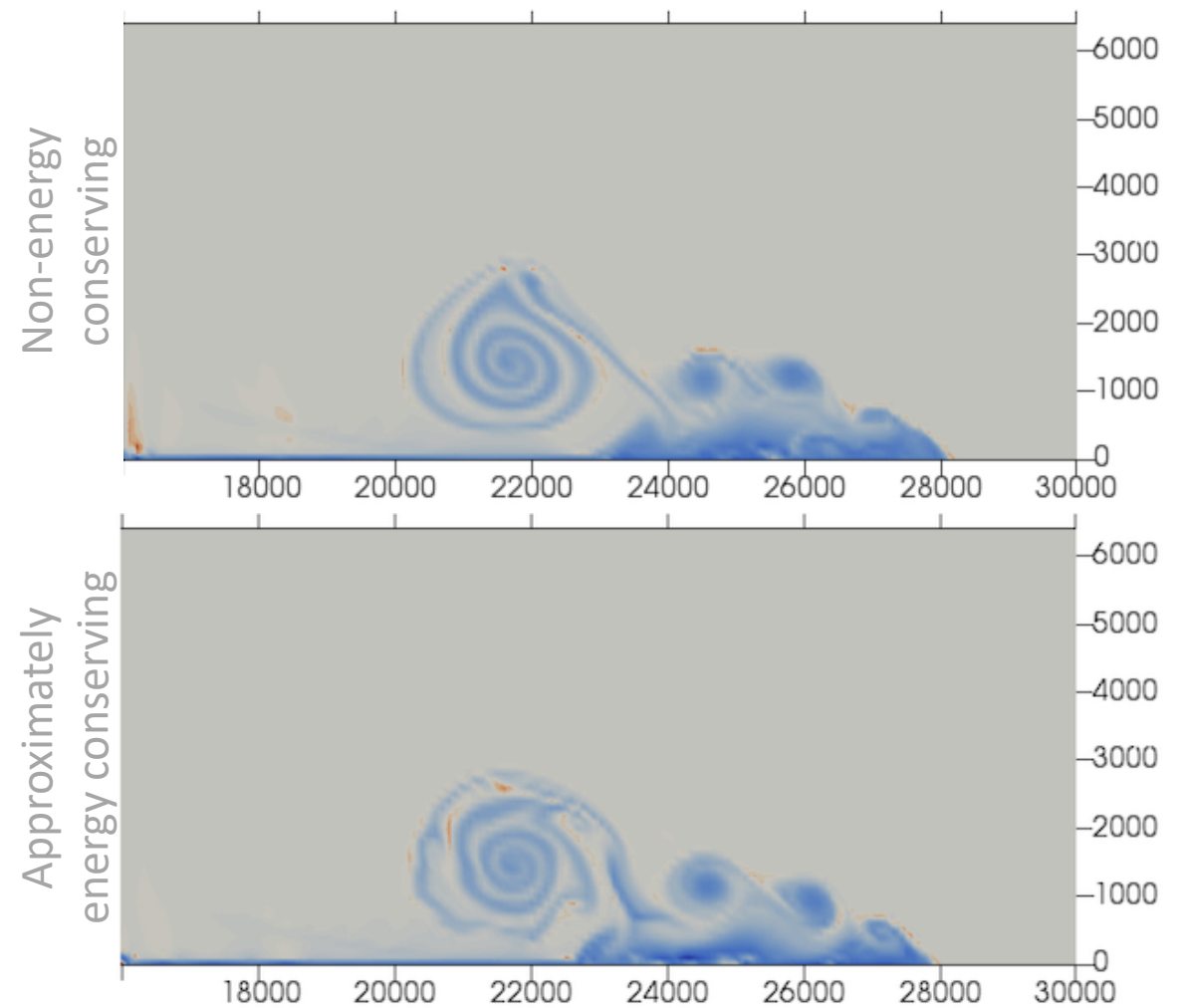
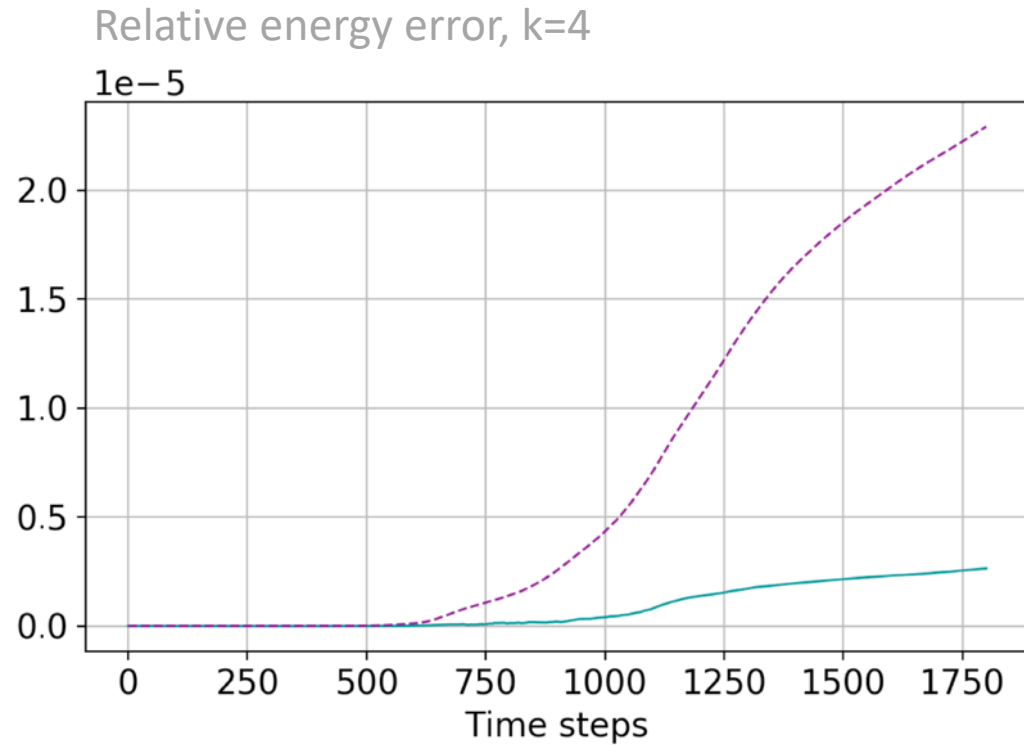
$$\overline{\frac{\delta H}{\delta \rho}} \in V_\rho \quad \overline{\frac{\delta H}{\delta \theta}} \in V_\theta$$

Relative energy error
relative to number of
quasi-Newton iterations



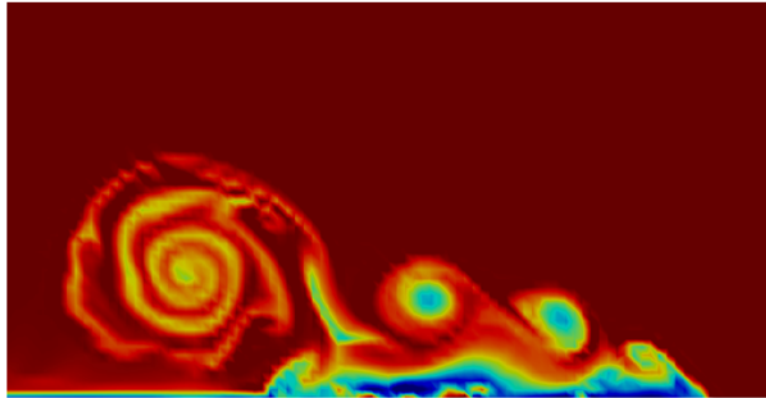
Results

- Falling cold air bubble

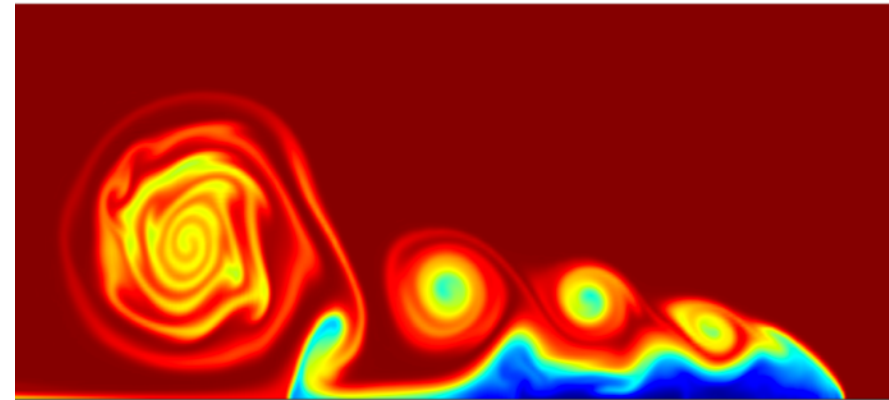


Results

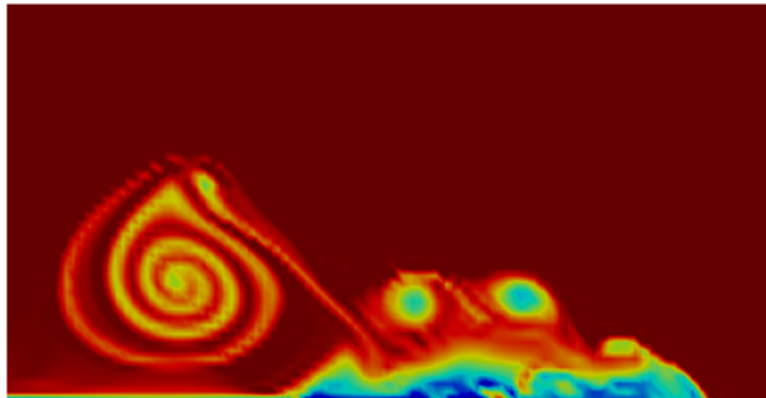
Approximately energy conserving



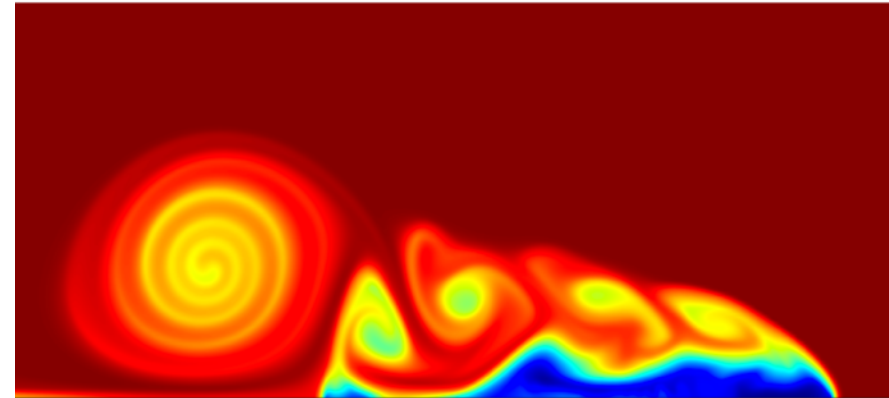
Dynamic sub-grid scale model



Non-energy conserving

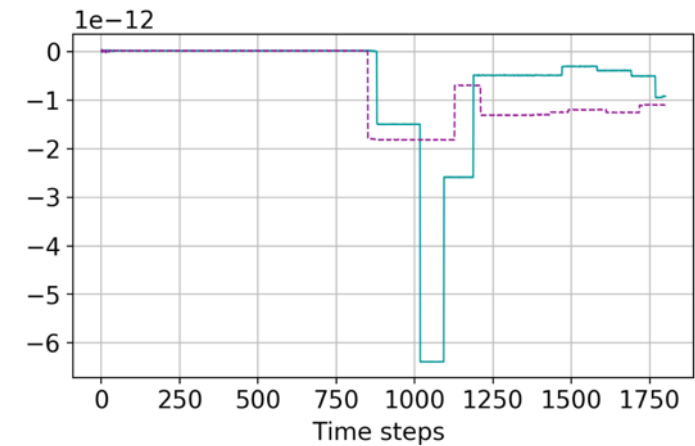
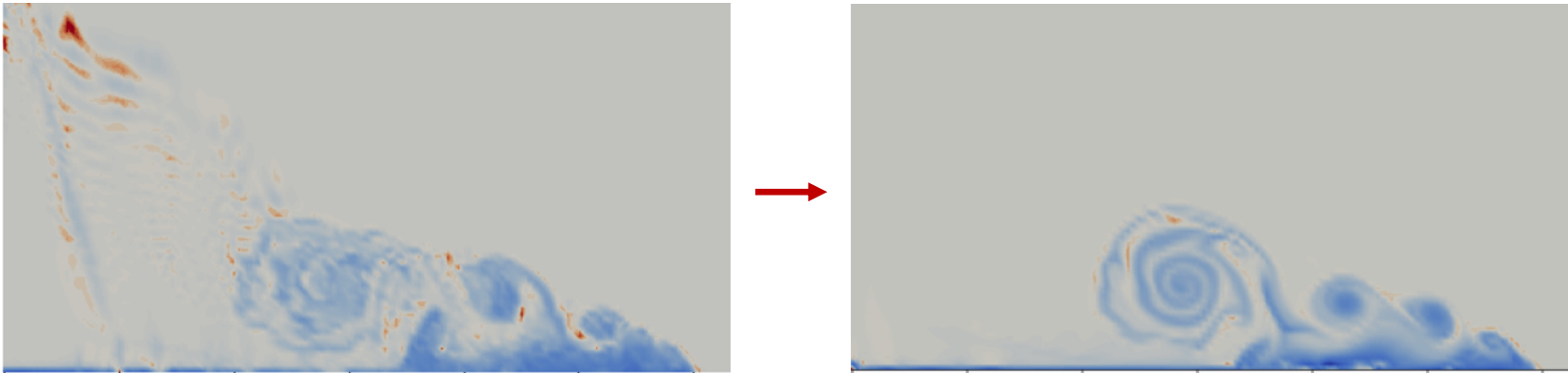


Smagorinsky sub grid scale model



Summary

- Fully upwinded, energy conserving formulation



- Simplification to reduce computational cost

$$\frac{\overline{\delta H}}{\delta \mathbf{u}} \longrightarrow \frac{\overline{\delta H}'}{\delta \mathbf{u}}$$

~~$$\mathbb{U}(\bar{\rho}, \frac{\overline{\delta H}}{\delta \mathbf{u}}), (\mathbb{S}_{\mathbf{u}}^{-1})^*, (\theta^a)^* \in V_u, \frac{\overline{\delta H}}{\delta \rho} \in V_\rho, \frac{\overline{\delta H}}{\delta \theta} \in V_\theta$$~~

References

- G.A. Wimmer, C.J. Cotter, W. Bauer. Energy conserving upwinded compatible finite element schemes for the rotating shallow water equations, *Journal of Computational Physics*, 401: 109016 (2020)
- G.A. Wimmer, C.J. Cotter, W. Bauer. Energy conserving SUPG methods for compatible finite element schemes in numerical weather prediction, *ArXiv preprint 2001.09590* (2020)
- S. Marras, M. Nazarov, F.X. Giraldo. Stabilized higher-order Galerkin methods based on a parameter-free dynamic SGS model for LES, *Journal of Computational Physics*, 301: 77-101 (2015)

