

1. Consider the system

$$\begin{aligned}\dot{x} &= \mu + x^2 + xy + y^2, \\ \dot{y} &= x^2 + xy - y, \\ \dot{\mu} &= 0.\end{aligned}$$

- Consider the critical point $(0, 0, 0)$ and compute the linearly stable, linearly unstable, and linearly neutral subspaces.
- Note that the system has a two-dimensional center manifold. Compute an expression for the center manifold up to terms of second order.
- Show that the motion on the center manifold is given by

$$\begin{aligned}\dot{u} &= \mu + u^2 + \text{h.o.t.}, \\ \dot{\mu} &= 0.\end{aligned}$$

- Now treat μ as a parameter and show that the system in center manifold coordinates has a bifurcation at $\mu = 0$. Name that bifurcation and sketch a bifurcation diagram in center manifold coordinates. (Extra credit if you manage to sketch the bifurcation in the original $2 + 1$ -dimensional phase space.)

(5+5+5+5)

(a) At $(x, y, \mu) = (0, 0, 0)$, the Jacobian matrix is

$$DF = \begin{pmatrix} 0 & 0 & 1 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

so that, by inspection, the system has a one-dimensional stable subspace $V_s = \text{span}\{e_2\}$ with eigenvalue -1 and a two-dimensional neutral gen. eigenspace $V_c = \text{span}\{e_1, e_3\}$.

(b) We can parameterize the center manifold W_c by x and μ .

Since it is tangent to the x - μ plane, the parametrization must start with quadratic terms, i.e.

$$y = ax^2 + bx\mu + c\mu^2 + \text{h.o.t.}$$

$$\Rightarrow \dot{y} = 2ax\dot{x} + b\dot{x}\mu + b x \dot{\mu} + 2c \mu \dot{\mu}$$

$$= (2ax + b\mu)(\mu + x^2 + xy + y^2)$$

on the other hand:

$$\dot{y} = x^2 + (x-1)(ax^2 + bx\mu + c\mu^2)$$

$$\Rightarrow 2ax\mu + b\mu^2 = x^2 - ax^2 - bx\mu - c\mu^2 + \text{h.o.t.}$$

$$\Rightarrow a=1, b=-2, c=2$$

(c) Insert back into the x -equation (with $x \equiv u$):

$$\dot{u} = \mu + u^2 + \text{h.o.t.}$$

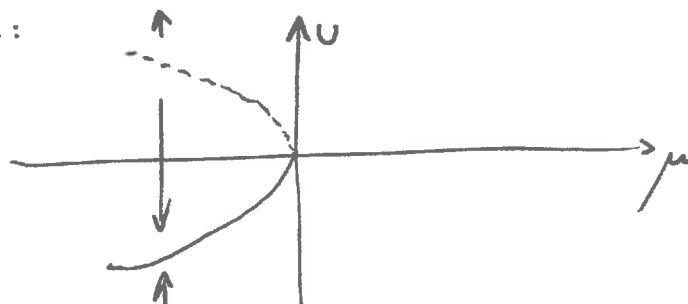
(The μ -equation remains unchanged)

(d) Critical points: $\mu + u^2 = 0$ (up to h.o.t.)

$$\Rightarrow u = \pm \sqrt{-\mu}$$

I.e. no critical point for $\mu > 0$, $u=0$ is crit. pt. for $\mu=0$, when $\mu < 0$, $u = \sqrt{-\mu}$ is unstable, $u = -\sqrt{-\mu}$ is stable.

It is a saddle-node:



2. Consider the system

$$\begin{aligned} \dot{x} &= \mu x + y - x(x^2 + y^2), \\ \dot{y} &= -x + \mu y - y(x^2 + y^2). \end{aligned}$$

follows most directly from (b) and (c).

- (a) Show that the origin is the only critical point and determine the linear stability of the origin as a function of the real parameter μ .
- (b) Write the system in polar coordinates.
- (c) Show that the system has a periodic orbit for $\mu > 0$. Conclude that the system has an Andronov-Hopf-bifurcation at $\mu = 0$.
- (d) Sketch the bifurcation diagram.

(5+5+5+5)

(a) At $(x, y) = (0, 0)$, the Jacobian matrix is

$$DF = \begin{pmatrix} \mu & 1 \\ -1 & \mu \end{pmatrix}$$

with eigenvalues given by $(\mu - \lambda)^2 + 1 = 0$

$$\Rightarrow \mu - \lambda = \pm i \Rightarrow \lambda = \mu \pm i$$

So for $\mu < 0$, the system is linearly stable, for $\mu = 0$, it is linearly neutrally stable, and for $\mu > 0$ it is linearly unstable.

(b) Write $x = r \cos \theta$, $y = -r \sin \theta$

$$\begin{aligned} \Rightarrow \dot{r} \cos \theta + r \dot{\theta} (-\sin \theta) &= \mu r \cos \theta - r \sin \theta - r^3 \cos \theta \\ -\dot{r} \sin \theta - r \dot{\theta} \cos \theta &= -r \cos \theta - \mu r \sin \theta + r^3 \sin \theta \end{aligned}$$

$$\begin{aligned} \Rightarrow \dot{r} &= \mu r - r^3 \\ -r \dot{\theta} &= -r, \text{ so } \dot{\theta} = 1 \end{aligned}$$

(c) Critical points of the τ -equation:

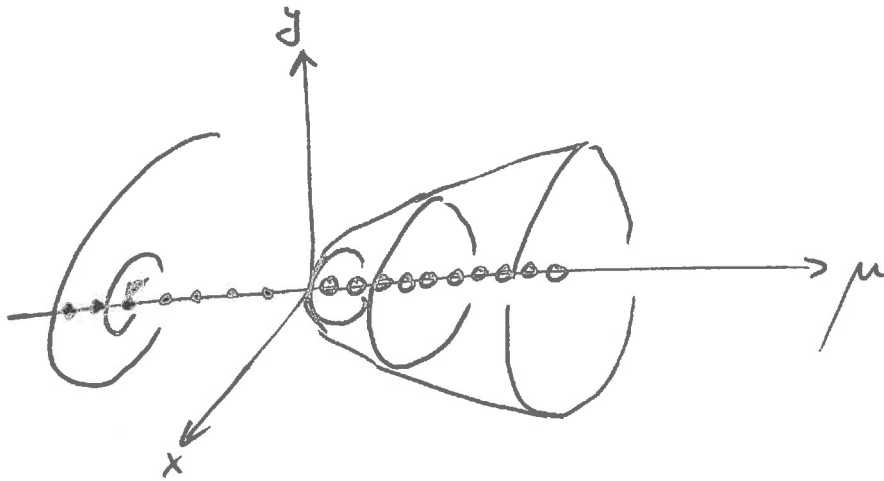
$$\mu\tau - \tau^3 = 0$$

$$\Rightarrow \tau = 0 \text{ or } \mu = \tau^2$$

When $\mu > 0$, there is a pair of critical points at $\tau = \pm\sqrt{\mu}$, clearly these are stable, since $\dot{\Theta} = 1$, it's a per. orbit in x - y variables

The critical point at $\tau = 0$ is stable for $\mu < 0$ and unstable for $\mu > 0 \Rightarrow$ Andronov-Hopf bifurcation at $\mu = 0$.

(d)



For $\mu < 0$, there is a stable focus, for $\mu > 0$ an unstable focus and a stable periodic orbit.

3. Prove the following simple averaging theorem.

Consider the equation

$$\dot{x} = \varepsilon f(x, t) + \varepsilon^2 g(x, t; \varepsilon),$$

where $f: \mathbb{R}^n \times [0, \infty) \rightarrow \mathbb{R}^n$ is T -periodic for fixed $x \in \mathbb{R}^n$. Assume that f and g are smooth. Suppose that there exists a ball of radius $R > 0$ on which f , g , and all their derivatives are bounded independent of ε .

Define

$$\bar{f}(y) = \frac{1}{T} \int_0^T f(y, t) dt$$

and suppose that the solution of

$$\dot{y} = \varepsilon \bar{f}(y) \quad y(0) = x(0)$$

remains in a ball of radius $R/2$. Then

$$x(t) - y(t) = O(\varepsilon)$$

on the time scale $1/\varepsilon$.

Hint: Define

$$u(x, t) = \int_0^t (f(x, s) - \bar{f}(x)) ds,$$

set

$$x = z + \varepsilon u(z, t)$$

and derive a differential equation for z .

(10)

Follow the hint:

$$\dot{x} = \dot{z} + \varepsilon \dot{u} = \dot{z} + \varepsilon (f(z, t) - \bar{f}(z)) + \varepsilon D_u \dot{z}$$

on the other hand,

$$\dot{x} = \varepsilon f(z + \varepsilon u, t) + \varepsilon^2 g$$

Note that, by definition, u is bounded uniformly in ε and t .

The same argument applies to D_u .

Thus, $x = z + O(\varepsilon)$ for as long as z remains in B_R .

Further, $g = O(\varepsilon^2)$ and $f(z + \varepsilon v, t) = f(z, t) + O(\varepsilon)$

for as long as x and z remain in B_R .

$$\Rightarrow (\mathbf{I} + \varepsilon D_v) \dot{z} = \varepsilon \bar{f}(z) + O(\varepsilon^2)$$

Under the assumptions above, and for ε small enough,

$\mathbf{I} + \varepsilon D_v$ is uniformly invertible with

$$(\mathbf{I} + \varepsilon D_v)^{-1} = \mathbf{I} - \varepsilon D_v + O(\varepsilon^2)$$

$$\Rightarrow \dot{z} = \varepsilon \bar{f}(z) + O(\varepsilon^2)$$

so that, by the Gronwall lemma, $z = y + O(\varepsilon)$ on the

time scale $\frac{1}{\varepsilon}$. We conclude that

$$x = y + O(\varepsilon)$$

on the time scale $\frac{1}{\varepsilon}$.

Note that the argument above is contingent on x and z

remaining in B_R . However, since y remains in $B_{\frac{R}{2}}$

by assumption and x, z remain $O(\varepsilon)$ -close, they cannot leave

B_R on the time scale $\frac{1}{\varepsilon}$ if ε is small enough.

4. Consider Mathieu's equation

$$\ddot{x} + (1 + 2\varepsilon \cos(2t))x = 0.$$

(a) Use the ansatz

$$x = z_1 \cos t + z_2 \sin t, \quad (*)$$

$$\dot{x} = -z_1 \sin t + z_2 \cos t \quad (**)$$

to show that

$$\left. \begin{aligned} \dot{z}_1 &= 2\varepsilon \sin t \cos(2t) (z_1 \cos t + z_2 \sin t), \\ \dot{z}_2 &= -2\varepsilon \cos t \cos(2t) (z_1 \cos t + z_2 \sin t). \end{aligned} \right\} (***)$$

(b) Use averaging (see attached table for useful trigonometric identities!) to conclude that

$$x(t) = \frac{1}{2} (x(0) + \dot{x}(0)) e^{-\frac{1}{2}\varepsilon t} (\cos t + \sin t) + \frac{1}{2} (x(0) - \dot{x}(0)) e^{\frac{1}{2}\varepsilon t} (\cos t - \sin t) + O(\varepsilon)$$

on the time scale $1/\varepsilon$. (Check the assumptions of the statement in Problem 3!)

(10+10)

(a) Differentiate (*) and use (**):

$$\dot{x} = \dot{z}_1 \cos t - z_1 \sin t + \dot{z}_2 \sin t + z_2 \cos t$$

$$\equiv -z_1 \sin t + z_2 \cos t$$

Differentiate (**) and use the equation:

$$\ddot{x} = -\dot{z}_1 \sin t - z_1 \cos t + \dot{z}_2 \cos t - z_2 \sin t$$

$$\equiv -(1 + 2\varepsilon \cos 2t) (z_1 \cos t + z_2 \sin t)$$

$$\Rightarrow \dot{z}_1 \cos t + \dot{z}_2 \sin t = 0$$

$$-\dot{z}_1 \sin t + \dot{z}_2 \cos t = -2\varepsilon \cos 2t (z_1 \cos t + z_2 \sin t)$$

$$\Rightarrow \dot{z}_1 = 2\varepsilon \sin t \cos 2t (z_1 \cos t + z_2 \sin t)$$

$$\dot{z}_2 = -2\varepsilon \cos t \cos 2t (z_1 \cos t + z_2 \sin t)$$

$$(b) \sin t \cos t \cos 2t = \frac{1}{2} \sin 2t \cos 2t = \frac{1}{4} \sin 4t$$

→ averages to 0

$$\begin{aligned} \sin^2 t \cos 2t &= \frac{1}{2} (1 - \cos 2t) \cos 2t \\ &= \frac{1}{2} \cos 2t - \frac{1}{4} (1 + \cos 4t) \end{aligned}$$

→ averages to $-\frac{1}{4}$

$$\begin{aligned} \cos^2 t \cos 2t &= \frac{1}{2} (1 + \cos 2t) \cos 2t \\ &= \frac{1}{2} \cos 2t + \frac{1}{4} (1 + \cos 4t) \end{aligned}$$

→ averages to $\frac{1}{4}$

⇒ the averaged system reads

$$\begin{aligned} \dot{y}_1 &= -\frac{1}{2}\varepsilon y_2 \\ \dot{y}_2 &= -\frac{1}{2}\varepsilon y_1 \end{aligned} \quad \text{or} \quad \dot{y} = \begin{pmatrix} 0 & -\frac{\varepsilon}{2} \\ -\frac{\varepsilon}{2} & 0 \end{pmatrix} y$$

This linear system has ⁹ eigenvalues $\lambda_1 = -\frac{\varepsilon}{2}$ and $\lambda_2 = \frac{\varepsilon}{2}$

with eigenvectors $v_1 = \begin{pmatrix} 1 \\ 1 \end{pmatrix}$ and $v_2 = \begin{pmatrix} 1 \\ -1 \end{pmatrix}$.

$$\Rightarrow y(t) = c_1 \begin{pmatrix} 1 \\ 1 \end{pmatrix} e^{-\frac{\epsilon}{2}t} + c_2 \begin{pmatrix} 1 \\ -1 \end{pmatrix} e^{\frac{\epsilon}{2}t}$$

$$\Rightarrow x(t) = (c_1 e^{-\frac{\epsilon}{2}t} + c_2 e^{\frac{\epsilon}{2}t}) \cos t + (c_1 e^{-\frac{\epsilon}{2}t} - c_2 e^{\frac{\epsilon}{2}t}) \sin t$$

$$\dot{x}(t) = -(c_1 e^{-\frac{\epsilon}{2}t} + c_2 e^{\frac{\epsilon}{2}t}) \sin t + (c_1 e^{-\frac{\epsilon}{2}t} - c_2 e^{\frac{\epsilon}{2}t}) \cos t$$

At $t=0$:

$$\begin{cases} x(0) = c_1 + c_2 \\ \dot{x}(0) = c_1 - c_2 \end{cases} \Rightarrow c_1 = \frac{x(0) + \dot{x}(0)}{2}, \quad c_2 = \frac{x(0) - \dot{x}(0)}{2}$$

Inserting these expressions into the solution for $x(t)$, we obtain the result as claimed.

It remains to verify that the assumptions of the averaging theorem are satisfied. We look at (**). Structure, periodicity, and smoothness match trivially. It remains to check boundedness of the averaged solution. Indeed, for $t = \frac{1}{\epsilon}$,

$$\|y(t)\| \leq |c_1| \sqrt{2} e^{-\frac{1}{2}} + |c_2| \sqrt{2} e^{\frac{1}{2}}$$

Choosing R twice this RHS, we can proceed.

5. Consider the system

$$\left. \begin{aligned} \varepsilon \dot{y} &= z + g(\phi), \\ \varepsilon \dot{z} &= -y, \\ \dot{\phi} &= I, \\ \dot{I} &= 0, \end{aligned} \right\} \Rightarrow -\varepsilon^2 \ddot{z} = z + g(\phi)$$

where $\phi \in \mathbb{T}$, where $\mathbb{T} = [0, 2\pi]$ with identification of the endpoints.

Show that the y - z dynamics can be parameterized via $z = z(\phi, I)$ provided the resonance condition

$$I \neq \frac{1}{\varepsilon k}, \quad k \in \mathbb{Z}$$

is satisfied.

Hint: Fix I , expand $g(\phi)$ and $z(\phi, I)$ as a Fourier series, and match coefficients. (10)

The angle ϕ is 2π -periodic, so we can write

$$g(\phi) = \sum_{k \in \mathbb{Z}} g_k e^{ik\phi}$$

$$z(\phi, I) = \sum_{k \in \mathbb{Z}} z_k e^{ik\phi} \quad \Rightarrow \quad \ddot{z} = \sum_{k \in \mathbb{Z}} z_k (-k^2 I^2) e^{ik\phi}$$

$$\Rightarrow \varepsilon^2 \sum_{k \in \mathbb{Z}} z_k k^2 I^2 e^{ik\phi} = \sum_{k \in \mathbb{Z}} z_k e^{ik\phi} + \sum_{k \in \mathbb{Z}} g_k e^{ik\phi}$$

$$\Rightarrow \varepsilon^2 k^2 I^2 z_k = z_k + g_k \quad \text{for every } k \in \mathbb{Z}$$

$$\Rightarrow g_k = \frac{z_k}{1 - \varepsilon^2 k^2 I^2}$$

10

So we can solve for g_k provided $1 - \varepsilon^2 k^2 I^2 \neq 0$,

i.e. if $I \neq \frac{1}{\varepsilon k}$ for any $k \in \mathbb{Z}$.