

## 2.6 Vector Operators

Session 13  
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Certain common operations involving partial derivatives have names.

As before, we use  $\vec{\nabla} = \begin{pmatrix} \frac{\partial}{\partial x_1} \\ \vdots \\ \frac{\partial}{\partial x_n} \end{pmatrix}$  as a vector that can operate on functions in different ways:

• For  $\varphi: \mathbb{R}^n \rightarrow \mathbb{R}$  (a scalar field),  $\vec{\nabla}\varphi = \begin{pmatrix} \frac{\partial \varphi}{\partial x_1} \\ \vdots \\ \frac{\partial \varphi}{\partial x_n} \end{pmatrix} =: \text{grad } \varphi$  is called **gradient of  $\varphi$** .

• For  $\vec{f}: \mathbb{R}^n \rightarrow \mathbb{R}^n$  (a vector field),  $\vec{\nabla} \cdot \vec{f} = \begin{pmatrix} \frac{\partial}{\partial x_1} \\ \vdots \\ \frac{\partial}{\partial x_n} \end{pmatrix} \cdot \begin{pmatrix} f_1 \\ \vdots \\ f_n \end{pmatrix} = \sum_{i=1}^n \frac{\partial f_i}{\partial x_i} =: \text{div } \vec{f}$   
is called **divergence of  $\vec{f}$** .

• For  $\vec{f}: \mathbb{R}^3 \rightarrow \mathbb{R}^3$  (a vector field),  $\vec{\nabla} \times \vec{f} = \begin{pmatrix} \frac{\partial}{\partial x_1} \\ \frac{\partial}{\partial x_2} \\ \frac{\partial}{\partial x_3} \end{pmatrix} \times \begin{pmatrix} f_1 \\ f_2 \\ f_3 \end{pmatrix} = \begin{pmatrix} \frac{\partial f_3}{\partial x_2} - \frac{\partial f_2}{\partial x_3} \\ \frac{\partial f_1}{\partial x_3} - \frac{\partial f_3}{\partial x_1} \\ \frac{\partial f_2}{\partial x_1} - \frac{\partial f_1}{\partial x_2} \end{pmatrix} =: \text{curl } \vec{f}$   
is called **curl of  $\vec{f}$** .

These operations get really interesting in the context of line/surface/volume integrals (Gauss and Stokes theorems), but we don't have the time to go into this.

Let us here just note some interesting identities:

$$\cdot \text{curl grad } \varphi = \vec{\nabla} \times (\vec{\nabla} \varphi) = \begin{pmatrix} \frac{\partial}{\partial x_1} \\ \frac{\partial}{\partial x_2} \\ \frac{\partial}{\partial x_3} \end{pmatrix} \times \begin{pmatrix} \frac{\partial \varphi}{\partial x_1} \\ \frac{\partial \varphi}{\partial x_2} \\ \frac{\partial \varphi}{\partial x_3} \end{pmatrix} = \begin{pmatrix} \frac{\partial x_2 \frac{\partial \varphi}{\partial x_3} - \frac{\partial x_3 \frac{\partial \varphi}{\partial x_2}}{\partial x_1} \\ -\frac{\partial x_1 \frac{\partial \varphi}{\partial x_3} + \frac{\partial x_3 \frac{\partial \varphi}{\partial x_1}}{\partial x_2} \\ \frac{\partial x_1 \frac{\partial \varphi}{\partial x_2} - \frac{\partial x_2 \frac{\partial \varphi}{\partial x_1}}{\partial x_3} \end{pmatrix} = 0$$

Clairaut/Schwarz

$$\cdot \text{div curl } \vec{f} = \begin{pmatrix} \frac{\partial}{\partial x_1} \\ \frac{\partial}{\partial x_2} \\ \frac{\partial}{\partial x_3} \end{pmatrix} \cdot \begin{pmatrix} \frac{\partial x_2 f_3 - \frac{\partial x_3 f_2}{\partial x_1}}{\partial x_1} \\ -\frac{\partial x_1 f_3 + \frac{\partial x_3 f_1}{\partial x_2}}{\partial x_2} \\ \frac{\partial x_1 f_2 - \frac{\partial x_2 f_1}{\partial x_3}}{\partial x_3} \end{pmatrix}$$

$$= \frac{\partial}{\partial x_1} (\frac{\partial x_2 f_3 - \frac{\partial x_3 f_2}{\partial x_1}}{\partial x_1}) + \frac{\partial}{\partial x_2} (-\frac{\partial x_1 f_3 + \frac{\partial x_3 f_1}{\partial x_2}}{\partial x_2}) + \frac{\partial}{\partial x_3} (\frac{\partial x_1 f_2 - \frac{\partial x_2 f_1}{\partial x_3}}{\partial x_3}) = 0$$

Clairaut/Schwarz

More examples in the homework/moodle exercises.

# 3. Ordinary Differential Equations

## 3.1 Basic Introduction

In many applications, we know relations between functions and their derivatives, e.g.,

- Newton's law:  $m \frac{d^2 \vec{x}(t)}{dt^2} = \vec{F}(\vec{x}(t))$ ; given  $\vec{F}$ , we want to find the particle trajectory  $\vec{x}(t)$   
e.g., Coulomb interaction  $\vec{F}(\vec{x}) = \text{const} \frac{\vec{x}}{|\vec{x}|^3}$   
(different possibilities for different initial positions and velocities)

- Also different derivatives (and complex numbers) might be involved, as, e.g., in the Schrödinger equation (in quantum mechanics):  
$$i \frac{\partial \psi(t,x)}{\partial t} = - \frac{\partial^2 \psi(t,x)}{\partial x^2} + V(x) \psi(t,x)$$
  
(1 dimensional equation for one particle;  $\psi: \mathbb{R}^2 \rightarrow \mathbb{C}$ )

- population growth:  $\frac{d\gamma(t)}{dt} = \lambda \gamma(t)$   
the more there is, the higher the increase  
Corona virus!  
( $\lambda > 0$  growth,  $\lambda < 0$  decay)

General setup:

Definition: For some given function  $f$ , we call

- $\underbrace{\gamma'(x)}_{= \frac{d\gamma(x)}{dx}} = f(x, \gamma(x))$  a first order ordinary differential equation (ODE)

( $\gamma: \mathbb{R} \rightarrow \mathbb{R}$  and  $f: \mathbb{R}^2 \rightarrow \mathbb{R}$  here).

If  $\gamma'(x) = f(\gamma(x))$  (no explicit  $x$ -dependence) we say the ODE is autonomous.

- $\underbrace{\gamma^{(n)}}_{n\text{-th derivative}}(x) = f(x, \gamma(x), \gamma'(x), \dots, \gamma^{(n-1)}(x))$  an  $n$ -th order ODE ( $\gamma: \mathbb{R} \rightarrow \mathbb{R}$  and  $f: \mathbb{R}^{n+1} \rightarrow \mathbb{R}$  here)

- $f(x_1, \dots, x_n, \gamma(\vec{x}), \frac{\partial \gamma}{\partial x_1}, \dots, \frac{\partial \gamma}{\partial x_n}, \frac{\partial^2 \gamma}{\partial x_1^2}, \frac{\partial^2 \gamma}{\partial x_1 \partial x_2}, \dots, \frac{\partial^n \gamma}{\partial x_1 \dots \partial x_n}) = 0$  an  $n$ -th order

all possible partial derivatives up to order  $n$

partial differential equation (PDE)

- Examples:
- population growth: 1st order ODE
  - Newton's law: 2nd order ODE
  - Schrödinger equation: 2nd order PDE

In this chapter we discuss some techniques to find solutions  $y(x)$  for certain types of equations.

(Only ODEs in this chapter.)

Ex.:  $\frac{dy}{dx} = \lambda y$ ,  $y: \mathbb{R} \rightarrow \mathbb{R}$ ,  $x \mapsto y(x)$ ,  $\lambda \in \mathbb{R}$  fixed

formally, we can bring all  $y$ 's and all  $x$ 's to different sides:  $\frac{dy}{y} = \lambda dx$ , and then

integrate:  $\int \frac{dy}{y} = \int \lambda dx \Rightarrow \ln y = \lambda x + C$

$\Rightarrow y(x) = e^{\lambda x + C} = e^C e^{\lambda x}$  is the solution!

this is called exponential growth for  $\lambda > 0$  (or decay for  $\lambda < 0$ )

We actually have the freedom to choose  $C$  here. How is  $C$  determined? By the value of  $y$  at any  $x_0$ :  $y(x_0) = e^C e^{\lambda x_0}$ , i.e., if  $x_0$  and  $y_0$  are given, we know  $C$ .

$\underbrace{y(x_0)}_{=: y_0}$

For some  $x_0$ , the  $y(x_0)$  is called initial condition.

More clearly, we could just write our solution as  $y(x) = \underbrace{y_0 e^{\lambda(x-x_0)}}_{\text{i.e., } e^C = y_0 e^{-\lambda x_0}}$  (s.t.  $y_0$  is initial condition at  $x_0$ ).

$\nearrow y(x_0) = y_0 e^{\lambda(x_0-x_0)} = y_0$

$\downarrow$

Often, one just chooses  $x_0 = 0$ , s.t.  $y_0 = y(0)$ .

Generally there is this important fact:

For an  $n$ -th order ODE  $y^{(n)}(x) = f(x, y(x), y'(x), \dots, y^{(n-1)}(x))$ , we need to specify  $n$  initial conditions  $y(x_0), y'(x_0), \dots, y^{(n-1)}(x_0)$ . In other words, the solution needs to have  $n$  independent constants.

Example: initial position  $x(t_0)$  and initial velocity  $x'(t_0)$  for Newton's equation  $\frac{d^2 \vec{x}(t)}{dt^2} = \vec{F}(\vec{x}(t))$

To summarize, the most important technique for solving ODEs is this:

### Separation of variables:

this is certainly not always possible, e.g.  $\frac{dy}{dx} = \cos(xy)$

For  $\frac{dy}{dx} = f(y, x)$ , we bring all  $x$ 's to one side and all  $y$ 's to the other (if possible), and then integrate both sides.

also here, we might or might not be able to actually perform the integration

## 3.2 Some Types of Integrable ODEs

integrable = find explicit solution by integration (Technique 1 above)

•  $y'(x) = f(x)g(y)$  is called separable ODE

here, we can write  $\frac{dy}{dx} = f(x)g(y)$ , i.e.,  $\frac{dy}{g(y)} = f(x)dx$

$\Rightarrow$  we find the solution by integrating,  $\int \frac{dy}{g(y)} = \int f(x)dx$  (if we can)

(note: at least we know that a solution exists if  $f$  and  $g$  are continuous and  $g(y) \neq 0 \forall y$ )  
because then we can integrate

•  $y'(x) = f(x)y(x)$  is called linear homogeneous ODE

as before:  $\frac{dy}{dx} = f(x)y \Rightarrow \frac{dy}{y} = f(x)dx \Rightarrow \int \frac{dy}{y} = \int f(x)dx$

$\Rightarrow \ln y = \int f(x)dx + C \Rightarrow y(x) = e^{\int f(x)dx + C}$  (so we can always find a solution as long as  $f$  can be integrated)

•  $y'(x) = f(x)y(x) + g(x)$  is called linear inhomogeneous ODE

here the idea is to write  $y(x) = \underbrace{u(x)v(x)}_{\text{some product of two fct.s}}$  (s.t. applying the product rule gives a sum of two functions)

$$\Rightarrow y'(x) = (u(x)v(x))' = \frac{du}{dx}v + u\frac{dv}{dx} = f(x)u(x)v(x) + g(x)$$

$$\Rightarrow \text{solve first } \frac{du}{dx} = f(x)u : \frac{du}{u} = f(x)dx \xrightarrow{\text{as before}} u(x) = e^{\int f(x)dx}$$

$$\text{next we need to solve } u\frac{dv}{dx} = g(x) \text{ i.e., } dv = \frac{g(x)}{u(x)} dx = e^{-\int f(x)dx} g(x) dx$$

$$\Rightarrow v(x) = \int \underbrace{e^{-\int f(x)dx}}_{\substack{\text{this notation means} \\ \text{this is a fct. of } x}} g(x) dx + C$$

$$\Rightarrow \text{our solution is } y(x) = e^{\int f(x)dx} \left( \int e^{-\int f(x)dx} g(x) dx + C \right)$$

There are many examples, let us just give one (more in the exercises):

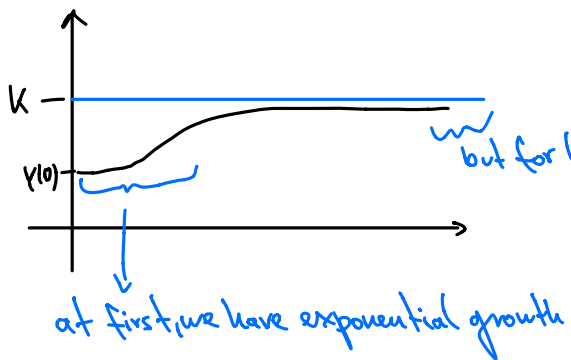
**Logistic growth:**  $\frac{dy}{dx} = \lambda y \left(1 - \frac{y}{k}\right)$  ( $\lambda$ : growth rate;  $k$  is sometimes called "environmental carrying capacity")

*this alone would lead to exponential growth*  $\rightarrow$  *growth is stopped once  $y$  reaches  $k$*

$\Rightarrow$  second order autonomous ODE

$$\text{separation of variables: } \frac{dy}{\lambda y \left(1 - \frac{y}{k}\right)} = dx$$

Integrating this will be a homework exercise. The result is:  $y(x) = e^{\lambda x} \left( C + \frac{e^{\lambda x}}{k} \right)^{-1}$



but for large  $x$ ,  $y(x) = \frac{e^{\lambda x}}{C + \frac{e^{\lambda x}}{K}} \approx \frac{e^{\lambda x}}{\frac{e^{\lambda x}}{K}} = K$