

Today we establish some standard properties of the Fourier transform on \mathcal{S} .

Lemma 2.11: Properties of the Fourier transform

$$(1) \quad \forall \alpha, \beta \in \mathbb{N}_0^d, f \in \mathcal{S}: \quad ((ik)^\alpha \partial_k^\beta \mathcal{F}f)(k) = (\mathcal{F} \partial_x^\alpha (-ix)^\beta f)(k),$$

in particular: $\widehat{(xf)}(k) = i(\widehat{\partial_k f})(k)$ and $\widehat{(\partial_x f)}(k) = ik\widehat{f}(k)$.

(2) \mathcal{F} and \mathcal{F}^{-1} are continuous linear maps $\mathcal{S} \rightarrow \mathcal{S}$.

Proof:

$$(1) \quad \text{recall } (\mathcal{F}f)(k) = \widehat{f}(k) = (2\pi)^{-\frac{d}{2}} \int f(x) e^{-ikx} dx$$

$$\text{Then with Lemma 2.2: } ((ik)^\alpha \partial_k^\beta \mathcal{F}f)(k) = (2\pi)^{-\frac{d}{2}} \int (ik)^\alpha \partial_k^\beta e^{-ikx} f(x) dx$$

note: all integrals exist
since $f \in \mathcal{S}$

$$\begin{aligned} &= (2\pi)^{-\frac{d}{2}} \int (ik)^\alpha (-ix)^\beta e^{-ikx} f(x) dx \\ &= (2\pi)^{-\frac{d}{2}} (-1)^{|\alpha|} \int (\partial_x^\alpha e^{-ikx}) (-ix)^\beta f(x) dx \end{aligned}$$

$|\alpha|$ -times integration by parts
(boundary terms vanish, since $f \in \mathcal{S}$)

$$\begin{aligned} &= (2\pi)^{-\frac{d}{2}} \int e^{-ikx} (\partial_x^\alpha (-ix)^\beta f(x)) dx \\ &= \mathcal{F}(\partial_x^\alpha (-ix)^\beta f)(k) \end{aligned}$$

(2) On metric spaces continuity (preimages of open sets are open) is equivalent to sequential continuity ($d(f_n, f) \rightarrow 0$ implies $d(\mathcal{F}f_n, \mathcal{F}f) \rightarrow 0$).

Thus let us choose $f_n \rightarrow f$ in S , meaning $d_S(f_n, f) \rightarrow 0$, meaning $\|f_n - f\|_{\alpha, \beta} \rightarrow 0$ for

all $\alpha, \beta \in \mathbb{N}_0^d$. We now show that $\|\mathcal{F}g\|_{\alpha, \beta} \leq C \sum_{j=0}^m \sup_{|\alpha|+|\beta|=j} \|g\|_{\tilde{\alpha}, \tilde{\beta}}$ for some $C > 0$ and $m \in \mathbb{N}$, which implies $\|\mathcal{F}f - \mathcal{F}f_n\|_{\alpha, \beta} \rightarrow 0 \quad \forall \alpha, \beta \in \mathbb{N}_0^d$ if $\|f_n - f\| \rightarrow 0$ $\forall \alpha, \beta \in \mathbb{N}_0^d$.

We compute:

$$\|\mathcal{F}g\|_{\alpha, \beta} = \|k^\alpha \partial_x^\beta \mathcal{F}g\|_\infty$$

$$\stackrel{(1) \text{ and } |\mathcal{S} \dots| \leq |\mathcal{S}| \dots}{\leq} (2\pi)^{-\frac{d}{2}} \int |\partial_x^\alpha \times^\beta g(x)| dx$$

$$= (2\pi)^{-\frac{d}{2}} \int (1+|x|^2)^d |\partial_x^\alpha \times^\beta g(x)| (1+|x|^2)^{-d} dx$$

$$\leq (2\pi)^{-\frac{d}{2}} \left(\sup_{x \in \mathbb{R}^d} (1+|x|^2)^d |\partial_x^\alpha \times^\beta g(x)| \right) \underbrace{\int (1+|x|^2)^{-d} dx}_{= \text{const.} \int_0^\infty (1+r^2)^{-d} r^{d-1} dr \leq \tilde{C}}$$

$$(\text{since integrand} \sim r^{-2d+d-1} \text{ for large } r)$$

$$\leq C \sum_{j=0}^m \sup_{|\alpha|+|\beta|=j} \|g\|_{\tilde{\alpha}, \tilde{\beta}} \quad \text{for some } m \in \mathbb{N}, C > 0.$$

□

Theorem 2.12:

$\mathcal{F}: S \rightarrow S$ is a continuous bijection with continuous inverse \mathcal{F}^{-1} .

Proof: HW: (1) Show $\mathcal{F}^{-1}\mathcal{F} = \text{id}$ only on $C_c^\infty = \text{smooth fct.s with compact support}$

↪ consider $\text{supp } f \subset [-m, m]^d$

⇒ Fourier series, write f as Riemann sum

(2) Show that C_c^∞ is dense in S , then thm. follows from continuity.

use some smooth cutoff function, e.g., $\delta(x) = \begin{cases} e^{-\frac{1}{1-x^2}+1} & \text{for } x < 1 \\ 0 & \text{else.} \end{cases}$

Lemma 2.14: Plancherel on S

For $f, g \in S$, we have $\int \hat{f}(x)g(x)dx = \int f(x)\hat{g}(x)dx$, and, in particular,

$$\int |\hat{f}(k)|^2 dk = \int |f(x)|^2 dx.$$

Proof: simple computation, HW.

Now back to the free SE $i\partial_t \Psi(t,x) = -\frac{1}{2} \Delta_x \Psi(t,x)$

Formally we solve this by applying \mathcal{F} : $i\partial_t \hat{\Psi}(t,k) = -\frac{1}{2} (\mathcal{F} \Delta_x \Psi)(t,k) \stackrel{\text{(Lemma 2.11)}}{=} \frac{1}{2} k^2 \hat{\Psi}(t,k)$

$$\Rightarrow \hat{\Psi}(t,k) = e^{-i\frac{k^2}{2}t} \hat{\Psi}(0,k) \text{ might global solution}$$

$$\Rightarrow \Psi(t,x) = (\mathcal{F}^{-1} e^{-i\frac{k^2}{2}t} \mathcal{F} \Psi_0)(x) \text{, with } \Psi_0(x) = \Psi(0,x) \text{ the initial condition}$$

As a map $\Psi: \mathbb{R} \rightarrow S$, this is indeed a differentiable solution to the free SE
(even ∞ -often differentiable). (Theorem next time.)