

4. Mean-field Dynamics for Bosons

4.1 Hartree Theory

We consider the Hamiltonian $H_N = \sum_{i=1}^N (-\Delta_i) + \frac{1}{N} \sum_{\substack{i,j=1 \\ i < j}}^N v(x_i - x_j)$ and the

associated many-body Schrödinger equation $i \frac{d}{dt} \Psi_N(t) = H_N \Psi_N(t)$, with $\Psi_N(t) \in L^2(\mathbb{R}^{3N})$

and $\Psi_N^t(x_1, \dots, x_N) = \Psi_N^t(x_{\sigma(1)}, \dots, x_{\sigma(N)}) \quad \forall \sigma \in S_N$ ($S_N =$ symmetric group = all permutations of $1, \dots, N$).

Ψ_N symmetric (bosons)

The choice of $\frac{1}{N}$ as coupling constant is called mean-field limit. It describes weak interaction, and is a simple model of a Bose-Einstein condensate (BEC).

Well-posedness:

With Kato-Rellich and HW 10, we find that H_N is self-adjoint on $\mathcal{D}(H_N) = H^2(\mathbb{R}^{3N})$

if $v: \mathbb{R}^3 \rightarrow \mathbb{R}$ even, and $v = v_1 + v_2$ with $v_1 \in L^2(\mathbb{R}^3)$, $v_2 \in L^\infty(\mathbb{R}^3)$ (we write $v \in L^2 + L^\infty$).

In this section, we aim at studying the dynamics of initial data $\Psi_N^{t=0}(x_1, \dots, x_N) = \prod_{i=1}^N \varphi^{t=0}(x_i)$, for some $\varphi^{t=0} \in L^2(\mathbb{R}^3)$.

most simple bosonic state

Such initial data mean that all particles are iid distributed.

We hope to prove that also $\Psi_N^t(x_1, \dots, x_N) \approx \prod_{i=1}^N \varphi^t(x_i)$ for some $\varphi(t) \in L^2(\mathbb{R}^3)$ in the limit $N \rightarrow \infty$.

What equation should hold for $\varphi(t)$?

- Idea: let X_i be an iid random variable with distribution $|\varphi^\dagger(x)|^2$.

Then $\frac{1}{N} \sum_{i=1}^N v(X_i - \gamma)$ should converge to $\int dx v(x - \gamma) |\varphi^\dagger(x)|^2 = (v * |\varphi^\dagger|^2)(\gamma)$ as $N \rightarrow \infty$ according to the law of large numbers.

- Thus we guess:
$$i \frac{\partial}{\partial t} \varphi(t, x) = -\Delta \varphi(t, x) + (v * |\varphi(t)|^2)(x) \varphi(t, x)$$
$$=: h^{\varphi(t)} \varphi(t, x)$$

This is called Hartree equation. It is a non-linear PDE, and one example of a non-linear SE (NLS). Thus, our previous well-posedness results for the linear SE do not apply. We come back to the question of well-posedness later.

- But note that formally:

$$\frac{d}{dt} \|\varphi(t)\|_{L^2}^2 = \frac{d}{dt} \langle \varphi(t), \varphi(t) \rangle = \langle \frac{d}{dt} \varphi(t), \varphi(t) \rangle + \langle \varphi(t), \frac{d}{dt} \varphi(t) \rangle$$

assuming $\varphi(t) \in H^2$ and $(v * |\varphi(t)|^2) \varphi(t) \in L^2$

$$\begin{aligned} &= i \langle (-\Delta \varphi(t) + (v * |\varphi(t)|^2) \varphi(t)), \varphi(t) \rangle - i \langle \varphi(t), (-\Delta \varphi(t) + (v * |\varphi(t)|^2) \varphi(t)) \rangle \\ &= 0 \quad (\text{integration by parts and } v * |\varphi(t)|^2 \in \mathbb{R}), \end{aligned}$$

So $\|\varphi(t)\|_{L^2} = \|\varphi(0)\|_{L^2}$ as for the linear SE

In order to prove $\Psi_N^\perp(x_1, \dots, x_N) \approx \prod_{i=1}^N \varphi^\perp(x_i)$, we proceed in several steps.

Step 1: Type of convergence

Definition 4.1:

For $\varphi \in L^2$, $\|\varphi\|_2 = 1$, we define the operator $p^\varphi: L^2(\mathbb{R}^3) \rightarrow L^2(\mathbb{R}^3)$, $\mathcal{X} \mapsto \langle \varphi, \mathcal{X} \rangle \varphi$, and

$q^\varphi := \mathbb{1} - p^\varphi$. For any $j=1, \dots, N$, we define $p_j^\varphi: L^2(\mathbb{R}^{3N}) \rightarrow L^2(\mathbb{R}^{3N})$ by

$(p_j^\varphi \Psi_N)(x_1, \dots, x_N) = \varphi(x_j) \int \overline{\varphi(y)} \Psi_N(x_1, \dots, x_{j-1}, y, x_{j+1}, \dots, x_N) dy$, and $q_j^\varphi := \mathbb{1} - p_j^\varphi$.

The following properties hold:

Lemma 4.2:

For any $\varphi \in L^2$ with $\|\varphi\|_2 = 1$, $j=1, \dots, N$ we have:

- (i) $p_j^\varphi q_j^\varphi \in \mathcal{L}(L^2(\mathbb{R}^{3N}))$ with $\|p_j^\varphi\|_{\mathcal{L}} = 1 = \|q_j^\varphi\|_{\mathcal{L}}$,
- (ii) $p_j^\varphi q_j^\varphi$ are orthogonal projectors ($P: \mathcal{H} \rightarrow \mathcal{H}$ is an orthogonal projector if $P^2 = P = P^*$),
- (iii) $p_j^\varphi q_k^\varphi = 0$, $[r_j^\varphi, s_k^\varphi] = 0$ for all j, k and $r, s \in \{p, q\}$.

Proof:

$$\begin{aligned}
 \text{(ii)} \quad \langle \mathcal{X}, p_j^\varphi \Psi \rangle &= \int dx_1 \dots dx_N \overline{\mathcal{X}(x_1, \dots, x_N)} \varphi(x_j) \int dy \overline{\varphi(y)} \Psi(x_1, \dots, y, \dots, x_N) \\
 &= \int dx_1 \dots dx_N dy \overline{\varphi(y)} \overline{\varphi(x_j)} \overline{\mathcal{X}(x_1, \dots, x_N)} \Psi(x_1, \dots, y, \dots, x_N) \\
 &= \int dx_1 \dots dy \dots dx_N \overline{\varphi(y)} \int dx_j \overline{\varphi(x_j)} \overline{\mathcal{X}(x_1, \dots, x_N)} \Psi(x_1, \dots, y, \dots, x_N) \\
 &= \langle p_j^\varphi \mathcal{X}, \Psi \rangle, \text{ so } p_j^\varphi = (p_j^\varphi)^*
 \end{aligned}$$

$$\begin{aligned}
\text{and } (p_j^q p_j^q \Psi)(x_1, \dots, x_n) &= \varphi(x_j) \int d\gamma \overline{\varphi(\gamma)} (p_j^q \Psi)(x_1, \dots, \gamma, \dots, x_n) \\
&= \varphi(x_j) \underbrace{\int d\gamma \overline{\varphi(\gamma)} \varphi(\gamma)}_{=\|\varphi\|_{L^2}^2=1} \int d\epsilon \overline{\varphi(\epsilon)} \Psi(x_1, \dots, \epsilon, \dots, x_n) \\
&= (p_j^q \Psi)(x_1, \dots, x_n).
\end{aligned}$$

$$\text{Also, } q_j^{q*} = \mathbb{1} - p_j^{q*} = \mathbb{1} - p_j^q = q_j^q \text{ and } q_j^{q^2} = (\mathbb{1} - p_j^q)(\mathbb{1} - p_j^q) = \mathbb{1} - 2p_j^q + p_j^{q^2} = \mathbb{1} - p_j^q = q_j^q.$$

$$(i) \ \|p_j^q \Psi\|_{L^2}^2 = \langle p_j^q \Psi, p_j^q \Psi \rangle = \langle \Psi, p_j^{q^2} \Psi \rangle = \langle \Psi, p_j^q \Psi \rangle \leq \|\Psi\| \|p_j^q \Psi\|$$

$$\Rightarrow \|p_j^q \Psi\|_{L^2} \leq \|\Psi\| \text{, i.e., } \|p_j^q\|_{\mathcal{L}} \leq 1$$

$$\text{Also: } p_j^q \prod_{i=1}^n \varphi(x_i) = \varphi(x_j) \int d\gamma \overline{\varphi(\gamma)} \varphi(x_1) \dots \varphi(\gamma) \dots \varphi(x_n) = \prod_{i=1}^n \varphi(x_i),$$

$$\text{So } \|p_j^q\|_{\mathcal{L}} := \sup_{\Psi, \|\Psi\|=1} \|p_j^q \Psi\| \geq \|p_j^q \prod_{i=1}^n \varphi(x_i)\| = \|\prod_{i=1}^n \varphi(x_i)\| = 1.$$

Same argument holds for q_j^q , with φ replaced by any $\varphi^\perp \in \{\varphi\}^\perp$.

$$(iii) \ p_j^q q_j^q = p_j^q (\mathbb{1} - p_j^q) = p_j^q - p_j^{q^2} = p_j^q - p_j^q = 0, \text{ and } r_j s_k = s_k r_j \text{ clear by def.} \quad \square$$

Note: $p_j^q \Psi_n$ tells us "how much" of the j -th particle is in the state q

$$\cdot \ p_j^q \prod_{i=1}^n \varphi(x_i) = \prod_{i=1}^n \varphi(x_i), \text{ and } q_j^q \prod_{i=1}^n \varphi(x_i) = 0$$