

We want to define multiplication operators $\psi(x) \mapsto f(x)\psi(x)$ as continuous maps on S , as we did with $e^{-i\frac{k^2}{2}t} \hat{\psi}_0$. For that, f cannot be too wild; an appropriate space is:

Definition 2.18: The space of smooth polynomially bounded functions is

$$C_{\text{pol}}^{\infty}(\mathbb{R}^d) := \left\{ f \in C^{\infty}(\mathbb{R}^d) : \forall \alpha \in \mathbb{N}_0^d \exists n_{\alpha} \in \mathbb{N} \text{ and } C_{\alpha} < \infty \text{ s.t. } |\partial_x^{\alpha} f(x)| \leq C_{\alpha} (1+|x|^2)^{\frac{n_{\alpha}}{2}} \right\}$$

Note: • a common notation is: $(1+|x|^2)^{\frac{1}{2}} =: \langle x \rangle$

• e.g., all polynomials $\in C_{\text{pol}}^{\infty}$, $e^{ikx} \in C_{\text{pol}}^{\infty}$, $e^x \notin C_{\text{pol}}^{\infty}$

Then indeed:

Lemma: For $f \in C_{\text{pol}}^{\infty}(\mathbb{R}^d)$, the multiplication operator $M_f: S \rightarrow S$, $\psi(x) \mapsto f(x)\psi(x)$ is continuous.

Proof: clear: if $\|\psi_n - \psi\|_{\alpha, \beta} \xrightarrow{n \rightarrow \infty} 0 \forall \alpha, \beta$, then also

$$\|M_f(\psi_n - \psi)\|_{\alpha, \beta} = \sup_{x \in \mathbb{R}^d} |x^{\alpha} \partial_x^{\beta} (f(x)(\psi_n(x) - \psi(x)))| \xrightarrow{n \rightarrow \infty} 0 \quad \square$$

The solution to the free SE can be written as $\mathcal{F}^{-1} e^{-i\frac{k^2}{2}t} \mathcal{F} \psi_0 = \mathcal{F}^{-1} M_f \mathcal{F} \psi_0$ for $f(k) = e^{-i\frac{k^2}{2}t}$. Since multiplication in Fourier space = derivatives in x -space, we introduce the following notation for $\mathcal{F}^{-1} M_f \mathcal{F}$:

Definition 2.19:

For $f \in C_{\text{pol}}^{\infty}(\mathbb{R}^d)$ we define the pseudo-differential operator

$$f(-i\mathcal{D}) : \mathcal{S}(\mathbb{R}^d) \rightarrow \mathcal{S}(\mathbb{R}^d), \quad \psi(x) \mapsto (f(-i\mathcal{D}_x)\psi)(x) = (\mathcal{F}^{-1} M_f \mathcal{F} \psi)(x) = (\mathcal{F}^{-1} f(k) \mathcal{F} \psi)(x)$$

Note: • $f(-i\mathcal{D})$ continuous, since $M_f, \mathcal{F}, \mathcal{F}^{-1}$ continuous

• $f(k) = k^{\alpha} \Rightarrow f(-i\mathcal{D}) = (-i)^{|\alpha|} \partial_x^{\alpha}$ is the usual differential operator

• Example: semi-relativistic or pseudo-relativistic Schrödinger equation:

$$i\partial_t \psi(t, x) = \underbrace{\sqrt{1 - \Delta}}_{\text{pseudo-differential operator}} \psi(t, x)$$

Examples:

• translation operator: for $a \in \mathbb{R}^d$, let $T_a(k) = e^{-iak} \Rightarrow T_a \in C_{\text{pol}}^{\infty}$

$$\begin{aligned} \Rightarrow \text{for } \psi \in \mathcal{S}, \text{ we find } (T_a(-i\mathcal{D})\psi)(x) &= (2\pi)^{-\frac{d}{2}} \int e^{ikx} e^{-iak} \hat{\psi}(k) dk \\ &= (2\pi)^{-\frac{d}{2}} \int e^{ik(x-a)} \hat{\psi}(k) dk \\ &= \psi(x-a) \end{aligned}$$

• free propagator: $P_f(k) = e^{-i\frac{k^2}{2}t} \Rightarrow P_f \in C_{\text{pol}}^{\infty}$

$$\begin{aligned} \Rightarrow \text{solution to free Schrödinger equation is } \psi(t, x) &= (\mathcal{F}^{-1} e^{-i\frac{k^2}{2}t} \mathcal{F} \psi_0)(x) \\ &= (P_f(t, -i\mathcal{D}) \psi_0)(x) \end{aligned}$$

Thus, we found: $\psi(t) = e^{-i\frac{(-\Delta)}{2}t} \psi(0)$

• heat equation: $\partial_t f(t,x) = \frac{1}{2} \Delta_x f(t,x)$

$$\Rightarrow W(t,k) = e^{-\frac{k^2}{2}t} \in C_{\text{pol}}^\infty \text{ for } t \geq 0$$

$$\Rightarrow \text{for } f(0,\cdot) = f_0 \in \mathcal{S}, t > 0, \text{ we have } f(t) = e^{\frac{1}{2}\Delta t} f_0 = W(t, -i\nabla) f_0$$

Lastly, $\mathcal{F}^{-1} M_f \mathcal{F} \psi_0 = \mathcal{F}^{-1} (f(x) \hat{\psi}_0(k))$, so we want to know about the (inverse) Fourier transform of a product.

Definition 2.22:

The convolution of $f \in \mathcal{S}$ and $g \in \mathcal{S}$ is $(f * g)(x) := \int_{\mathbb{R}^d} f(x-y) g(y) dy$.

Lemma 2.23: For $f, g, h \in \mathcal{S}$ we have

a) $(f * g) * h = f * (g * h)$ and $f * g = g * f$;

b) the map $\mathcal{S} \rightarrow \mathcal{S}, g \mapsto f * g$ is continuous;

c) $\widehat{f * g} = (2\pi)^{\frac{d}{2}} \hat{f} \cdot \hat{g}$ and $\widehat{f \cdot g} = (2\pi)^{-\frac{d}{2}} \hat{f} * \hat{g}$,

in particular $g(-i\nabla) f = \mathcal{F}^{-1} M_g \mathcal{F} f = \mathcal{F}^{-1} g \hat{f} = (2\pi)^{-\frac{d}{2}} \hat{g} * \hat{f}$.

Proof: • a) and c) are direct calculations

• then b) follows since $f * g = (2\pi)^{\frac{d}{2}} \mathcal{F}^{-1} \hat{f} \hat{g}$, i.e., composition of continuous maps \square

Example: heat equation: $f(t,x) = W(t, -i\nabla) f(0,x)$, $W(t,k) = e^{-\frac{k^2}{2}t}$
 $= (2\pi)^{-\frac{d}{2}} ((\mathcal{F}^{-1} W_t) * f_0)(x)$

with heat kernel $G(t,x) = (2\pi)^{-\frac{d}{2}} (\mathcal{F}^{-1} W)(t,x) = (2\pi t)^{-\frac{d}{2}} e^{-\frac{x^2}{2t}}$ we find

$$f(t, x) = (G(t) * f_0)(x) = (2\pi t)^{-\frac{d}{2}} \int e^{-\frac{(x-y)^2}{2t}} f_0(y) dy$$

To summarize: The solution to the free SE is

$$\Psi_0(t) = \mathcal{F}^{-1} M_{\mathcal{P}_f} \mathcal{F} \Psi_0 = e^{-i(\frac{\Delta}{2})t} \Psi_0 = G(t) * \Psi_0, \text{ with } \mathcal{P}_f(k) = e^{-i\frac{k^2}{2}t}, G(t, x) = (2\pi t)^{-\frac{d}{2}} e^{\frac{i(x-y)^2}{2t}}$$