

Next: define $\mathcal{F}, \mathcal{F}^{-1}, M_f : S' \rightarrow S'$

We start with \mathcal{F} :

Definition 2.31: $\mathcal{F}_{S'} := \mathcal{F}'_S$, meaning for $T \in S'$, we define its Fourier transform $\hat{T} \in S'$ by $\hat{T}(f) = T(\hat{f}) \quad \forall f \in S$.

Corollary 2.32: $\mathcal{F}' : S' \rightarrow S'$ is a weak*-continuous bijection, and $\hat{T}_f = T_{\hat{f}}$ for all $f \in S$ (or $f \in L^1$) (recall $T_f(g) := \int f g$).

i.e., $\hat{T}_f(g) = T_{\hat{f}}(g) = \int \hat{f} g \quad \forall g \in S$

Proof: $\mathcal{F} : S \rightarrow S$ is continuous and linear, so we conclude with Thm. 2.30 that $\mathcal{F}' : S' \rightarrow S'$ is weak*-continuous.

Bijective? $(\mathcal{F}'^{-1} \mathcal{F}' T)(f) = (\mathcal{F}' T)(\mathcal{F}'^{-1} f) = T(\mathcal{F} \mathcal{F}'^{-1} f) = T(f)$
 \Rightarrow yes, with continuous inverse $\mathcal{F}'^{-1} = \mathcal{F}^{-1}$.

Also, for $f \in S$ or $f \in L^1$:

$$\hat{T}_f(g) = (\mathcal{F}' T_f)(g) = T_f(\mathcal{F}' g) = \int f(x) \hat{g}(x) dx \stackrel{\text{Plancherel}}{=} \int \hat{f}(x) g(x) = T_{\hat{f}}(g) \quad \forall g \in S \quad \square$$

Ex.: Fourier transform of δ ($\delta(f) = f(0)$)

$$\Rightarrow \hat{\delta}(f) = \delta(\hat{f}) = \hat{f}(0) = \int \underbrace{(2\pi)^{-\frac{d}{2}}}_{g(x)} f(x) dx = T_g(f)$$

$\Rightarrow T_g$ with $g(x) = (2\pi)^{-\frac{d}{2}}$ is the Fourier transform of δ , or " $\hat{\delta}(k) = (2\pi)^{-\frac{d}{2}}$ "

Next: derivatives

Note: $\partial_x^\alpha : S \rightarrow S$ is linear (clear) and continuous, since

$$\|\partial_x^\alpha f\|_{S_1, \beta} = \|x^\beta \partial_x^\alpha f\|_\infty = \|f\|_{S_1, \alpha+\beta} \quad (\text{i.e., continuity on } S \text{ follows as usual from sequential continuity})$$

Definition 2.34: $\tilde{\partial}_x^\alpha := (-1)^{|\alpha|} \partial_x^\alpha : S' \rightarrow S'$, i.e., for $T \in S'$ the distributional derivative $\tilde{\partial}_x^\alpha T$ is defined by $(\tilde{\partial}_x^\alpha T)(f) := T((-1)^{|\alpha|} \partial_x^\alpha f) \quad \forall f \in S$.

Corollary 2.35: $\tilde{\partial}_x^\alpha : S' \rightarrow S'$ is weak*-continuous and $\tilde{\partial}_x^\alpha T_g = T_{\partial_x^\alpha g} \quad \forall g \in S$.

Proof: Weak*-continuity follows again from Thm. 2.30.

$$\begin{aligned} \text{Also, } (\tilde{\partial}_x^\alpha T_g)(f) &= T_g((-1)^{|\alpha|} \partial_x^\alpha f) = \int g(x) (-1)^{|\alpha|} \partial_x^\alpha f(x) dx \\ &\stackrel{\substack{\text{1st times} \\ \text{integration by} \\ \text{parts}}}{=} \int (\partial_x^\alpha g(x)) f(x) dx = T_{\partial_x^\alpha g}(f) \quad \forall f \in S. \end{aligned}$$

Ex.: • For $\theta(x) = \mathbb{1}_{[0, \infty)}(x) := \begin{cases} 1 & \text{for } x \geq 0 \\ 0 & \text{for } x < 0 \end{cases}$, we find $\frac{d}{dx} \theta = \delta$, see HW.

• $\tilde{\partial}_x^\alpha \delta$? See HW.

Summary: we have defined $\tilde{\mathcal{F}} = \mathcal{F}' : S' \rightarrow S'$ and $\tilde{\partial}_x^\alpha = (-1)^{|\alpha|} \partial_x^\alpha : S' \rightarrow S'$.

Furthermore one can show:

- Fixing $h \in \mathcal{S}$, we can define the convolution $h \tilde{*} \cdot : \mathcal{S}' \rightarrow \mathcal{S}'$ via $(h \tilde{*} T)(f) = T(\tilde{h} * f)$ with $\tilde{h}(x) = h(-x)$. This definition is chosen such that $h \tilde{*} T_g = T_{g * h}$ for $g \in \mathcal{S}$.

$$\begin{aligned} \rightarrow (h \tilde{*} T_g)(f) &:= T_g(\tilde{h} * f) = \int dx g(x) \int dy h(y-x) f(y) \\ &= \int dy f(y) \int dx h(y-x) g(x) \end{aligned}$$

- Fixing $g \in C_{\text{pol}}^{\infty}$, we define $\tilde{M}_g = M_g'$ i.e., $(M_g T)(f) = T(M_g f)$.

↳ Note: gT well-defined for $g \in C_{\text{pol}}^{\infty}$, but product of distributions a-priori undefined (much research effort to define it at least for some distributions, e.g., Heisenberg's regularity structures).

Both are weak* continuous maps.

Note: $\{T_f \in \mathcal{S}' : f \in \mathcal{S}\}$ is dense in \mathcal{S}' w.r.t. weak* topology (not obvious, proof omitted).

Thus, T_f allows us to identify \mathcal{S} with some subset of \mathcal{S}' .

Because of density and continuity of the adjoint, the definition $A' T_f = T_{A f}$ uniquely defines A' on all of \mathcal{S}' . ← This is why we defined, e.g., $\tilde{\partial}_x^{\alpha} = (-1)^{|\alpha|} \partial_x^{\alpha}$.

From now on, we will forget about $\tilde{\cdot}$ or $'$ in the notation for the adjoint.

⇒ We have defined $\mathcal{F}T, \partial_x^{\alpha} T, h * T$ for $h \in \mathcal{S}$, gT for $g \in C_{\text{pol}}^{\infty}$ ($T \in \mathcal{S}'$).

With that we can solve the free Schrödinger equation on \mathcal{S}' :

Theorem 2.40:

Let $\psi_0 \in S'$, then the unique global solution to the free Schrödinger equation

$$i\partial_t \psi = -\frac{1}{2}\Delta \psi \text{ (in the sense of distributions) with } \psi(0) = \psi_0 \text{ is } \psi(t) = \mathcal{F}^{-1} e^{-i\frac{k^2}{2}t} \mathcal{F} \psi_0,$$

with $\psi \in C^\infty(\mathbb{R}_t, S'(\mathbb{R}^d))$.

Proof: First, note that $\psi(t) = \mathcal{F}^{-1} e^{-i\frac{k^2}{2}t} \mathcal{F} \psi_0 \in S'$ since $\mathcal{F}, \mathcal{F}^{-1}, M_f: S' \rightarrow S'$.

Next, let us check if this $\psi(t)$ solves the SE. For any $f \in S$, we find

$$i \frac{d}{dt} (f, \psi(t))_{S, S'} = i \frac{d}{dt} (f, \mathcal{F}^{-1} e^{-i\frac{k^2}{2}t} \mathcal{F} \psi_0)_{S, S'}$$

$$\text{by def. } \rightarrow = i \frac{d}{dt} (\mathcal{F} e^{-i\frac{k^2}{2}t} \mathcal{F}^{-1} f, \psi_0)_{S, S'}$$

$$\text{continuity of } \psi_0: S \rightarrow \mathbb{C} \rightarrow = (\mathcal{F} \left(\frac{d}{dt} e^{-i\frac{k^2}{2}t} \right) \mathcal{F}^{-1} f, \psi_0)_{S, S'}$$

$$= (\mathcal{F} e^{-i\frac{k^2}{2}t} \underbrace{\frac{k^2}{2} \mathcal{F}^{-1} f}_{= \mathcal{F}^{-1}(-\frac{\Delta}{2} f)}, \psi_0)_{S, S'}$$

$$= \left(-\frac{\Delta}{2} f, \mathcal{F}^{-1} e^{-i\frac{k^2}{2}t} \mathcal{F} \psi_0 \right)_{S, S'}$$

$$\text{by def. of the } \rightarrow = (f, -\frac{\Delta}{2} \psi(t))_{S, S'}. \\ \text{distributional derivative}$$

Similarly $\left(i \frac{d}{dt} \right)^k (f, \psi(t))_{S, S'} = \left(\left(-\frac{\Delta}{2} \right)^k f, \psi(t) \right)_{S, S'}$, so $\psi(t) \in C^\infty(\mathbb{R}_t, S'(\mathbb{R}^d))$. \square