

Next we will prove that indeed  $\alpha(\psi_n(t), \varrho(t)) \xrightarrow{n \rightarrow \infty} 0 \quad \forall t \in \mathbb{R}$ .

Step 2: Controlling  $\alpha(\psi_n(t), \varrho(t))$

A standard technique is based on (variations of) the following lemma:

Lemma 4.11: Gronwall Lemma

Let  $\eta: \mathbb{R} \rightarrow \mathbb{R}$  be differentiable and satisfy  $\frac{d}{dt} \eta(t) \leq C(\eta(t) + \varepsilon)$  for some  $C, \varepsilon \geq 0$ .

Then, for all  $t > 0$ , we have

$$\eta(t) \leq e^{Ct} \eta(0) + (e^{Ct} - 1) \varepsilon$$

Proof: First, consider differentiable  $f: \mathbb{R} \rightarrow \mathbb{R}$  satisfying  $\frac{d}{dt} f(t) \leq C f(t)$ .

Let  $g(t) := e^{Ct}$  (i.e.,  $g(t) > 0$ ). Then

$$\frac{d}{dt} \left( \frac{f(t)}{g(t)} \right) = \frac{\frac{df(t)}{dt} g(t) - f(t) \frac{dg(t)}{dt}}{g(t)^2} \leq \frac{C f(t) g(t) - f(t) C g(t)}{g(t)^2} = 0.$$

$$\Rightarrow \frac{f(t)}{g(t)} \leq \frac{f(0)}{g(0)} \Rightarrow f(t) \leq \underbrace{g(t)}_{=e^{Ct}} f(0) \underbrace{\frac{1}{g(0)}}_{=1} = e^{Ct} f(0). \quad (*)$$

Next, define  $h(t) := e^{Ct} \eta(0) + (e^{Ct} - 1) \varepsilon$ , s.t.  $\frac{dh(t)}{dt} = C e^{Ct} \eta(0) + C e^{Ct} \varepsilon = C(h(t) + \varepsilon)$

and  $h(0) = \eta(0)$ .

Then  $\frac{d}{dt} (\eta(t) - h(t)) \leq C(\eta(t) + \varepsilon) - C(h(t) + \varepsilon) = C(\eta(t) - h(t))$ , so (\*) implies

$$\eta(t) - h(t) \leq e^{Ct} (\eta(0) - h(0)) = 0 \quad \text{i.e., } \eta(t) \leq h(t). \quad \square$$

We hope to apply the Gronwall lemma to  $\alpha(\psi_n(t), \varrho(t))$ .

So let us compute:

$$\frac{d}{dt} \alpha(\psi_\nu(t), \varrho(t)) = \frac{d}{dt} \langle \psi_\nu(t), q_1^{\varrho(t)} \psi_\nu(t) \rangle$$

$$= \underbrace{\langle \frac{d}{dt} \psi_\nu(t), q_1^{\varrho(t)} \psi_\nu(t) \rangle}_{= -i H_\nu \psi_\nu(t) \text{ for } \psi_\nu(0) \in H^2(\mathbb{R}^{2n})} + \langle \psi_\nu(t), q_1^{\varrho(t)} \frac{d}{dt} \psi_\nu(t) \rangle + \langle \psi_\nu(t), \underbrace{\left( \frac{d}{dt} q_1^{\varrho(t)} \right) \psi_\nu(t)}_{= -\frac{d}{dt} p_1 = -\frac{d}{dt} |\varrho(t)| \langle \varrho(t) |} \rangle$$

If Heisenberg eq. holds in  $L^2$  sense, which we assume here.

$$\begin{aligned} &= - \left( -i h^{\varrho(t)} \varrho(t) \langle \varrho(t) | + |\varrho(t) \rangle \langle -i h^{\varrho(t)} \varrho(t) | \right) \\ &= - \left( -i h^{\varrho(t)} p + i p^{\varrho(t)} h^{\varrho(t)} \right) \\ &= i [h^{\varrho(t)}, p^{\varrho(t)}] \\ &= -i [h_1^{\varrho(t)}, q_1^{\varrho(t)}] \end{aligned}$$

$$\Rightarrow \frac{d}{dt} \alpha(\psi_\nu(t), \varrho(t)) = i \langle H_\nu \psi_\nu(t), q_1^{\varrho(t)} \psi_\nu(t) \rangle - i \langle \psi_\nu(t), q_1^{\varrho(t)} H_\nu \psi_\nu(t) \rangle - i \langle \psi_\nu(t), [h_1^{\varrho(t)}, q_1^{\varrho(t)}] \psi_\nu(t) \rangle$$

$$= i \langle \psi_\nu(t), [H_\nu - h_1^{\varrho(t)}, q_1^{\varrho(t)}] \psi_\nu(t) \rangle$$

$$= \left[ \sum_{j=1}^N (-\Delta_j) + \frac{1}{N-1} \sum_{\substack{i,j=1 \\ i \neq j}}^N v_{ij} - (-\Delta_1) - (v * |\varrho(t)|^2)_1, q_1^{\varrho(t)} \right]$$

$=: v(x_i - x_j)$

$$\begin{aligned} [-\Delta_j, q_1^{\varrho(t)}] &= 0 \text{ for } j > 1 \\ [v_{ij}, q_1^{\varrho(t)}] &= 0 \text{ for } i \neq 1 \text{ and } j \neq 1 \end{aligned}$$

$$= \left[ \frac{1}{N-1} \sum_{j=2}^N v_{1j} - (v * |\varrho(t)|^2)_1, q_1^{\varrho(t)} \right]$$

$$\Rightarrow \frac{d}{dt} \alpha(\psi_\nu(t), \varrho(t)) = i \langle \psi_\nu(t), \left[ \frac{1}{N-1} \sum_{j=2}^N v_{1j} - (v * |\varrho(t)|^2)_1, q_1^{\varrho(t)} \right] \psi_\nu(t) \rangle$$

$\psi_\nu(t)$  symmetric  $\Rightarrow i \langle \psi_\nu(t), [v_{12} - (v * |\varrho(t)|^2)_1, q_1^{\varrho(t)}] \psi_\nu(t) \rangle$

For  $A, B$  symmetric:  $\langle \psi, [A, B] \psi \rangle = -2 \operatorname{Im} \langle \psi, (v_{12} - (v * |\varrho(t)|^2)_1) q_1^{\varrho(t)} \psi \rangle$

$$= \langle A\psi, B\psi \rangle - \langle B\psi, A\psi \rangle = \langle A\psi, B\psi \rangle - \overline{\langle A\psi, B\psi \rangle} = 2i \operatorname{Im} \langle A\psi, B\psi \rangle$$

$$= -2 \operatorname{Im} \langle \psi, p_1^{\varrho(t)} \underbrace{(v_{12} - (v * |\varrho(t)|^2)_1)}_{=: W_{12}} q_1^{\varrho(t)} \psi \rangle$$

term with  $q_1^{\varrho(t)}$  here vanishes since then the scalar product is real

Inserting also  $\mathbb{1} = p_2^{\varphi(t)} + q_2^{\varphi(t)}$  leads to

$$\frac{d}{dt} \alpha(\psi_\nu(t), \varphi(t)) = -2 \operatorname{Im} \langle \psi_\nu(t), p_1^{\varphi(t)} p_2^{\varphi(t)} W_{12} q_1^{\varphi(t)} p_2^{\varphi(t)} \psi_\nu(t) \rangle \quad \text{term (I)}$$

$$-2 \operatorname{Im} \langle \psi_\nu(t), p_1^{\varphi(t)} q_2^{\varphi(t)} W_{12} q_1^{\varphi(t)} p_2^{\varphi(t)} \psi_\nu(t) \rangle \quad \text{term (II)}$$

$\in \mathbb{R}$  (since  $W_{12} = W_{21}$ )

$$-2 \operatorname{Im} \langle \psi_\nu(t), p_1^{\varphi(t)} p_2^{\varphi(t)} W_{12} q_1^{\varphi(t)} q_2^{\varphi(t)} \psi_\nu(t) \rangle \quad \text{term (III)}$$

$$-2 \operatorname{Im} \langle \psi_\nu(t), p_1^{\varphi(t)} q_2^{\varphi(t)} W_{12} q_1^{\varphi(t)} q_2^{\varphi(t)} \psi_\nu(t) \rangle \quad \text{term (IV)}$$