

Last time we proved:

$$\text{Let } \alpha(\Psi_\nu(t), \varrho(t)) := \langle \Psi_\nu(t), q_1^{\varrho(t)} \Psi_\nu(t) \rangle, \text{ with } i \frac{d}{dt} \Psi_\nu(t) = H_\nu \Psi_\nu(t), \quad i \frac{d}{dt} \varrho(t) = h^{\varrho(t)} \varrho(t),$$

$$\text{where } H_\nu = -\sum_{j=1}^N (-\Delta_j) + \frac{1}{N-1} \sum_{i,j} v(x_i - x_j), \quad h^{\varrho(t)} = -\Delta + v * |\varrho(t)|^2, \quad q_1^{\varrho(t)} := 1 - p^{\varrho(t)}, \quad p^{\varrho(t)} = |\varrho(t)\rangle \langle \varrho(t)|.$$

Let $\Psi_\nu(0) \in H^2(\mathbb{R}^{3N})$, $\varrho(t) \in H^2(\mathbb{R}^3)$. Then

$$\frac{d}{dt} \alpha(\Psi_\nu(t), \varrho(t)) = -2 \operatorname{Im} \langle \Psi_\nu(t), p_1^{\varrho(t)} p_2^{\varrho(t)} W_{12} q_1^{\varrho(t)} q_2^{\varrho(t)} \Psi_\nu(t) \rangle \quad \text{term (I)}$$

$$-2 \operatorname{Im} \underbrace{\langle \Psi_\nu(t), p_1^{\varrho(t)} q_2^{\varrho(t)} W_{12} q_1^{\varrho(t)} p_2^{\varrho(t)} \Psi_\nu(t) \rangle}_{\in \mathbb{R} \text{ (since } W_{12} = W_{21})}} \quad \text{term (II)} = 0$$

$$-2 \operatorname{Im} \langle \Psi_\nu(t), p_1^{\varrho(t)} p_2^{\varrho(t)} W_{12} q_1^{\varrho(t)} q_2^{\varrho(t)} \Psi_\nu(t) \rangle \quad \text{term (III)}$$

$$-2 \operatorname{Im} \langle \Psi_\nu(t), p_1^{\varrho(t)} q_2^{\varrho(t)} W_{12} q_1^{\varrho(t)} q_2^{\varrho(t)} \Psi_\nu(t) \rangle \quad \text{term (IV)}$$

$$\text{where } W_{12} := v_{12} - \underbrace{(v * |\varrho(t)|^2)_1}_{(v * |\varrho(t)|^2)_1}$$

$$\begin{aligned} \text{Now note that } p_2^{\varrho(t)} W_{12} p_2^{\varrho(t)} &= p_2^{\varrho(t)} v_{12} p_2^{\varrho(t)} - \underbrace{(v * |\varrho(t)|^2)_1}_{(v * |\varrho(t)|^2)_1} p_2^{\varrho(t)} = 0. \\ &= |\varrho(t)\rangle_2 \langle \varrho(t)|_2 v_{12} |\varrho(t)\rangle_2 \langle \varrho(t)|_2 \\ &= (v * |\varrho(t)|^2)_1 \\ &= (v * |\varrho(t)|^2)_1 p_2^{\varrho(t)} \end{aligned}$$

Thus, term(I) = 0. This is the essential term where the interaction v_{12} is cancelled by its average $v * |\varrho|^2$.

Furthermore: $|\text{term (III)}| \leq 2 \underbrace{\|p_1^{(t)}\|_{q_2} \|p_2^{(t)}\|_{q_1} \|\Psi_N(t)\|}_{\leq \|\Psi_N(t)\|} \underbrace{\|W_{12}\|_{\mathcal{S}}}_{\leq \|v\|_{L^\infty} + \|v * |\varphi(t)|^2\|_{L^\infty}} \underbrace{\|q_1^{(t)}\|_{q_2} \|q_2^{(t)}\|_{q_1} \|\Psi_N(t)\|}_{\leq \sqrt{\alpha(\Psi_N(t), \varphi(t))}}$.

$\leq \|\Psi_N(t)\|$
 $= \sqrt{\alpha(\Psi_N(t), \varphi(t))}$

$\leq \|v\|_{L^\infty} + \|v * |\varphi(t)|^2\|_{L^\infty} \leq \sqrt{\alpha(\Psi_N(t), \varphi(t))}$
 $\leq \|v\|_{L^\infty} \|\varphi(t)\|_{L^2}^2 = 1$
 Similar to HW3

$\Rightarrow |\text{term (III)}| \leq 4 \|v\|_{L^\infty} \alpha(\Psi_N(t), \varphi(t)).$

Finally: $\text{term (II)} = -2 \langle \Psi_N(t), p_1^{(t)} p_2^{(t)} (v_{12} - (v * |\varphi(t)|^2)_1) q_1^{(t)} q_2^{(t)} \Psi_N(t) \rangle$

$= -2 \langle \Psi_N(t), p_1^{(t)} p_2^{(t)} v_{12} q_1^{(t)} q_2^{(t)} \Psi_N(t) \rangle$

$\underbrace{p_2^{(t)} (v * |\varphi(t)|^2)_1 q_2^{(t)}}_{=0}$

$\underbrace{p_2^{(t)} q_2^{(t)} (v * |\varphi(t)|^2)_1}_{=0}$

We prove in HW 11 that $|\text{term (II)}| \leq 6 \|v\|_{L^\infty} (\alpha + \frac{1}{N})$ (for all $N \geq 3$).

Using Gronwall's lemma, we have proven:

Theorem 4.12: Derivation of the Hartree equation

Assume v is even and $v \in L^\infty$. Let $\Psi_N(t)$ be the solution to the Schrödinger equation with symmetric initial data $\Psi_N(0) \in H^2(\mathbb{R}^{3N}), \|\Psi_N(0)\| = 1$. Assume the Hartree equation has a unique solution $\varphi(t)$ given initial data $\varphi(0) \in L^2(\mathbb{R}^3), \|\varphi(0)\| = 1$.

We assume $\varphi(t) \in H^2(\mathbb{R}^3)$. Then

$$\alpha(\Psi_N(t), \varphi(t)) \leq e^{Ct} \alpha(\Psi_N(0), \varphi(0)) + (e^{Ct} - 1) \frac{1}{N} \quad \text{for } C = 10 \|v\|_{L^\infty}.$$

Note: The assumption that $\varphi(t) \in H^2(\mathbb{R}^3)$ exists holds for any $\varphi(0) \in H^2(\mathbb{R}^3)$. This can be proven with a Gronwall argument as well, but we skip that here.

Note: We can relax the condition $v \in L^\infty$. Let us consider $v(x) = |x|^{-1}$. Then:

- H_ν is still self-adjoint by Kato-Rellich.

- Term (I) = 0 still.

- $| \text{Term (III)} | \leq 2 \underbrace{\| q_2^{(t)}(\psi_w(t)) \|}_{\leq \sqrt{\alpha(t)}} \| p_1^{(t)} W_{12} \|_{\mathcal{L}} \underbrace{\| q_1^{(t)} \|_{\mathcal{L}}}_{\leq 1} \underbrace{\| q_2^{(t)}(\psi_w(t)) \|}_{\leq \sqrt{\alpha(t)}}$

and $\| p_1^{(t)} W_{12} \|_{\mathcal{L}} = \| W_{12} p_1^{(t)} \|_{\mathcal{L}} \leq \| v_{12} p_1^{(t)} \|_{\mathcal{L}} + \| v^* |q(t)|^2 \|_{\mathcal{L}} \| p_1^{(t)} \|_{\mathcal{L}}$.

Now: $\| v_{12} p_1^{(t)} \|_{\mathcal{L}}^2 \leq \| p_1^{(t)} v_{12} p_1^{(t)} \|_{\mathcal{L}} \leq \| v^* |q(t)|^2 \|_{\mathcal{L}} = \sup_{\gamma \in \mathbb{R}^3} \int dx v(x) |q(t, \gamma-x)|^2$
 $= (v^* |q(t)|^2)_2 p_1^{(t)}$

$= \sup_{\gamma \in \mathbb{R}^3} \langle \varphi_\gamma(t), \frac{1}{|x|^2} \varphi_\gamma(t) \rangle$, with $\varphi_\gamma(t, x) = \varphi(t, \gamma-x)$

$\leq 4 \langle \varphi(t), (-\Delta) \varphi(t) \rangle$ by Hardy's inequality (HW5 Problem 6)

$= 4 \int dx \varphi(t, \gamma-x) (-\Delta_x) \varphi(t, \gamma-x)$

$= 4 \int dx \varphi(t, x) (-\Delta_x) \varphi(t, x)$

$= 4 \| \nabla \varphi(t) \|^2$

Also: $\| v^* |q(t)|^2 \|_{\mathcal{L}} \leq \| v^2 |q(t)|^2 \|_{\mathcal{L}} \leq 4 \| \nabla \varphi(t) \|^2$

$\Rightarrow \| p_1^{(t)} W_{12} \|_{\mathcal{L}} \leq 8 \| \nabla \varphi(t) \|^2$

- $| \text{Term (II)} | \leq C \| \nabla \varphi(t) \|^2 (\alpha(t) + \frac{1}{\nu})$ can be shown as well.

- Now recall energy conservation from HW12 Problem 2:

For $E(\varphi(t)) := \| \nabla \varphi(t) \|^2 + \frac{1}{2} \int (v^* |q(t)|^2)(x) |\varphi(t, x)|^2 dx$, we have $E(\varphi(t)) = E(\varphi(0))$.

$\Rightarrow \| \nabla \varphi(t) \|^2 = \underbrace{E(\varphi(t))}_{= E(\varphi(0))} - \frac{1}{2} \int (v^* |q(t)|^2)(x) |\varphi(t, x)|^2 dx \leq E(\varphi(0))$.

Conclusion: Derivation of Hartree equation still works if $v(x) = \frac{1}{|x|}$ (or, in fact, $v = v^{(1)} + v^{(2)}$ with $v^{(1)} \in L^\infty$ and $|v^{(2)}(x)| \leq \frac{C}{|x-2|}$).